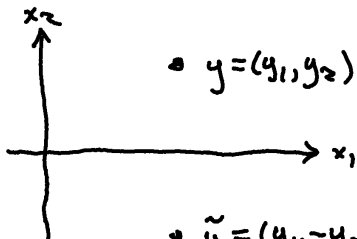


# TMA430S 2020: Øving for uke 47

B12.3 Husk fundamentalløsninga i  $\mathbb{R}^2$ :

(a)

$$\Phi(x) = \frac{-\ln r}{2\pi} = -\frac{\ln \sqrt{x_1^2 + x_2^2}}{2\pi} = -\frac{1}{4\pi} \ln(x_1^2 + x_2^2)$$



•  $y = (y_1, y_2)$

•  $\tilde{y} = (y_1, -y_2)$

Green's funksjon i  $\mathbb{H}^2$  bli-

$$G_y(x) = \underbrace{\Phi(y-x) - \Phi(\tilde{y}-x)}_{\text{harmonisk i } \mathbb{H}^2}$$

lik  $\Phi(y-x)$

når  $x_2 = 0$

skrevet fullt ut:

$$G_y(x) = -\frac{1}{4\pi} \left( \ln((x_1 - y_1)^2 + (x_2 - y_2)^2) - \ln((x_1 - y_1)^2 + (x_2 + y_2)^2) \right)$$

Løsningen av  $-\Delta u = 0$  i  $\mathbb{H}^2$ ,  $u(x, 0) = g(x)$  vil bli:

$$u(y) = -\int_{\partial \mathbb{H}^2} \partial_{\nu} G_y g \, dS^1 = \int_{\mathbb{R}} \partial_{x_2} G_y(x, 0) g(x) \, dx$$

(her:  $\partial_{\nu} = -\partial_{x_2}$ )

$$\partial_{x_2} G_y(x, 0) = -\frac{1}{4\pi} \frac{-4y_2}{(x-y_1)^2 + y_2^2} = \frac{1}{\pi} \frac{y_2}{(x-y_1)^2 + y_2^2} = P_{y_2}(x-y_1)$$

der  $P_y(x) = \frac{1}{\pi} \frac{y}{x^2 + y^2}$  er Poisson-kjernen for  $\mathbb{H}^2$ .

(heldt analogt med Poisson-kjernen for  $\mathbb{D}$ )

$$\begin{aligned} \text{Merk at } P_y(x) > 0, \text{ og } \int_{\mathbb{R}} P_y(x) \, dx &= \frac{y}{\pi} \int_{\mathbb{R}} \frac{dx}{x^2 + y^2} \\ &= \frac{y^2}{\pi} \int_{\mathbb{R}} \frac{dt}{(t^2 + 1)y^2} = \frac{1}{\pi} \int_{\mathbb{R}} \frac{dt}{1+t^2} = 1 \end{aligned}$$

$$\begin{aligned} x &= yt \\ dx &= y dt \end{aligned}$$

(b) Løsningsformelen blir  $u(y) = P_{y_2} * g(y_1)$

Fordi  $P_y(x) = \frac{1}{y} P_1\left(\frac{x}{y}\right)$  og  $\int_{\mathbb{R}} P_1 dx = 1$ , er

$\lim_{y \downarrow 0} P_y = \delta$  i  $\mathcal{D}'(\mathbb{R})$ , slik at  $\lim_{y_2 \downarrow 0} u(y) = \int_{\mathbb{R}} \delta * g(y_1) = g(y_1)$

(I hvert fall for  $g \in C_c^\infty(\mathbb{R})$ , men det er lett nok  
å utvide dette til  $g \in C_b(\mathbb{R}) := (C(\mathbb{R}) \cap L^\infty(\mathbb{R}))$ .)

10.1 Jeg skriver like gjerne  $xu'(x) = 1$ .

(a) I distributionsforstand:  $(xu', \varphi) = (1, \varphi)$

kan skrives  $(u', x\varphi) = (1, \varphi)$

altså  $-(u, (x\varphi)') = (1, \varphi)$

Som integral:  $-\int_{\mathbb{R}} u(x)(x\varphi(x))' dx = \int_{\mathbb{R}} \varphi dx$

Om  $u \in \mathcal{D}'$  og  $\varphi \in C^1_0$   
så er  $\varphi u \in \mathcal{D}'$   
definerent vel  
 $(\varphi u, \varphi) = (u, \varphi \varphi)$

(b)  $u = \ln|x|$  er lokalt integral fordi det ubestemte  
integral er  $x \ln|x| - x$  som er kontinuerlig i 0  
(om du setter verdien i 0 like).

$-\int_{\mathbb{R}} \ln|x| (x\varphi(x))' dx$  kan deles i to:

$$-\int_{-\infty}^0 \ln|x| (x\varphi(x))' dx = -\lim_{\varepsilon \rightarrow 0} \ln \varepsilon \cdot (-\varepsilon \varphi(\varepsilon)) + \int_{-\infty}^0 \frac{1}{x} \cdot x\varphi(x) dx$$

$$\stackrel{=0}{\text{for}} \varepsilon \ln \varepsilon \rightarrow 0 \quad = \int_{-\infty}^0 \varphi(x) dx$$

og  $\varphi$  er begrenset

$$-\int_0^{\infty} \ln|x| (x\varphi(x))' dx = \int_0^{\infty} \varphi(x) dx \text{ er samme grunn.}$$

$$\text{Så } -\int_{\mathbb{R}} \ln|x| \cdot (x\varphi(x))' dx = \int_{\mathbb{R}} \varphi(x) dx.$$

10.7  $u \in C^2(\bar{\Omega})$  og  $f \in C(\bar{\Omega})$ ; anta

$$\int_{\Omega} (\nabla u \cdot \nabla \psi - f\psi) d^4x = 0 \quad \text{for alle } \psi \in C^1(\bar{\Omega}).$$

Fordi  $\nabla \cdot (\psi \nabla u) = \psi \Delta u + \nabla \psi \cdot \nabla u$ , or

$$\int_{\Omega} (\psi \Delta u + \nabla u \cdot \nabla \psi) d^4x = \int_{\partial \Omega} \psi \partial_{\nu} u dS \quad (\text{Gauss's l. identitet})$$

Kombinerer man de to gr-

$$\int_{\Omega} (-\Delta u - f)\psi d^4x = -\int_{\partial \Omega} \psi \partial_{\nu} u dS'$$

Spesielt or  $\int_{\Omega} (-\Delta u - f)\psi d^4x = 0$  for  $\psi \in C_c^{\infty}(\Omega)$ ,

si  $-\Delta u = f$ . Men si blir  $\int_{\partial \Omega} \psi \partial_{\nu} u dS' = 0$

for alle  $\psi \in C_c^{\infty}(\mathbb{R}^n)$ , si  $\partial_{\nu} u = 0$ .