## Harmonicfunctionology

Harald Hanche-Olsen

The Laplace operator on  $\mathbb{R}^n$ 

$$\Delta = \sum_{i=1}^{n} \frac{\partial^2}{\partial x_i^2}$$

plays a rôle in the wave and heat equations, but even more fundamentally, in the *Laplace equation* 

$$\Delta u = 0 \tag{1}$$

and the *Poisson equation*  $-\Delta u = f$  where *f* is a given function.

A  $C^2$  solution of (1) is called *harmonic*. (Later, we will find that harmonic functions are in fact infinitely differentiable.)

Much is gained by the trivial observation that  $\Delta u = \nabla \cdot \nabla u$  together with various applications of the divergence theorem or two of its corollaries, Green's first and second identities.

Let's get started by simply integrating  $\Delta u$  over a bounded domain  $\omega$  with piecewise  $C^1$  boundary:

$$\int_{\omega} \Delta u(\mathbf{x}) \,\mathrm{d}^{n} \mathbf{x} = \int_{\omega} \nabla \cdot \nabla u(\mathbf{x}) \,\mathrm{d}^{n} \mathbf{x} = \int_{\partial \omega} \partial_{\nu} u(\mathbf{x}) \,\mathrm{d}S(\mathbf{x}). \tag{2}$$

This immediately proves

**Proposition 1.** If a  $C^2$  function u on a domain  $\Omega$  is harmonic, then

$$\int_{\partial \omega} \partial_{\nu} u \, \mathrm{d}S = 0 \tag{3}$$

for all bounded domains  $\omega \subseteq \Omega$  with piecewise  $C^1$  boundary.

Conversely, if (3) holds for every ball  $\omega = B(\mathbf{x}, r)$  whose closure lies within  $\Omega$ , then u is harmonic.

*Proof.* We have already proved the first part. For the converse, (3) and (2) imply that the *average* of  $\Delta u$  over any ball is zero. By letting the radius of the ball  $B(\mathbf{x}, r)$  tend to zero, we conclude that  $\Delta u(\mathbf{x}) = 0$ .

**Definition.** The (*radius r*) *spherical average* of a function u at a point x is defined to be

$$\tilde{u}_{\mathbf{x}}(r) = \int_{\partial B(\mathbf{x},r)} u \, \mathrm{d}S = \int_{\mathbb{S}^{n-1}} u(\mathbf{x} + r\mathbf{y}) \, \mathrm{d}S(\mathbf{y}),$$

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where  $\mathbb{S}^{n-1} \subset \mathbb{R}^n$  is the unit sphere and the "barred" integral signs denote the average:

$$\int_{\partial B(x,r)} u \, \mathrm{d}S = \frac{1}{A_n r^{n-1}} \int_{\partial B(x,r)} u \, \mathrm{d}S,$$

where  $A_n$  is the area of  $\mathbb{S}^{n-1}$ . Note that the second integral in the definition of  $\tilde{u}_x(r)$  makes sense even for r < 0; thus, we adopt this as the definition for all real r for which the integrand is defined on  $\mathbb{S}^{n-1}$ . We see that  $\tilde{u}_x$  is an *even* function; it is  $C^k$  if u is  $C^k$ , and  $\tilde{u}_x(0) = u(x)$ .

When  $\omega$  is a ball, we can rewrite (2) in terms of spherical averages: Noting that the volume of the ball  $B(\mathbf{x}, r)$  is  $A_n r^n / n$ , we find

$$\int_{B(\boldsymbol{x},r)} \Delta u(\boldsymbol{y}) \, \mathrm{d}^{n} \boldsymbol{y} = \frac{n}{r} \int_{B(\boldsymbol{x},r)} \partial_{\nu} u(\boldsymbol{y}) \, \mathrm{d}S(\boldsymbol{y}) = \frac{n}{r} \int_{\mathbb{S}^{n-1}} \partial_{r} u(\boldsymbol{x}+r\boldsymbol{y}) \, \mathrm{d}S(\boldsymbol{y}),$$

where we can move the r derivative outside the integral, and arrive at

$$\int_{B(\mathbf{x},r)} \Delta u(\mathbf{y}) \,\mathrm{d}^n \mathbf{y} = \frac{n}{r} \tilde{u}'_{\mathbf{x}}(r). \tag{4}$$

Along with  $\tilde{u}_{\mathbf{x}}(0) = u(\mathbf{x})$ , this implies

**Theorem 2** (The mean value property of harmonic functions). A  $C^2$  function u on a domain  $\Omega$  is harmonic if and only if  $\tilde{u}_x(r) = u(\mathbf{x})$  for all  $x \in \Omega$  and all r for which  $\overline{B}(\mathbf{x}, |\mathbf{r}|) \subset \Omega$ .

In general, we say a function *u* satisfies the *mean value property* if  $\tilde{u}_{\mathbf{x}}(r) = u(\mathbf{x})$  whenever  $\overline{B}(\mathbf{x}, |r|) \subset \Omega$ . We hall see below (Theorem 4) that the mean value property characterizes harmonic functions. But first, we collect an easy consequence of (4).

**Proposition 3.** For any  $C^2$  function u, we have

$$\Delta u(\mathbf{x}) = n\tilde{u}_{\mathbf{x}}''(0).$$

*Proof.* The function  $\tilde{u}_x$  is even, so  $\tilde{u}'_x(0) = 0$ . Therefore, letting  $r \to 0$  in (4), we arrive at the stated result.

**Theorem 4** (The mean value property and regularity). Assume that a continuous function u satisfies the mean value property on a domain  $\Omega$ . Then u is infinitely differentiable, and is therefore harmonic. In particular, every harmonic function is infinitely differentiable.

*Proof.* This proof may seem long, but only because we use it to develop some tools that have wider applicability.

First, define a *standard mollifier*  $\rho : \mathbb{R}^n \to \mathbb{R}$ . Here is one of many possible definitions:

$$\rho(\mathbf{x}) = \begin{cases} ae^{1/(|\mathbf{x}|^2 - 1)}, & |\mathbf{x}| < 1\\ 0, & |\mathbf{x}| \ge 1, \end{cases}$$

where a > 0 is chosen to ensure that

$$\int_{\mathbb{R}^n} \rho \, \mathrm{d} \boldsymbol{x} = 1.$$

That is one of the defining qualities of a standard mollifier. The others are: That  $\rho \ge 0$  everywhere, that it vanishes outside the unit ball, that it is infinitely differentiable, and is radially symmetric – that is, a function of  $|\mathbf{x}|$  alone.

For any  $\delta > 0$  we can squeeze the mollifier to fit inside a ball of radius  $\delta$ :

$$\rho_{\delta}(\boldsymbol{x}) = \frac{1}{\delta^n} \rho\left(\frac{\boldsymbol{x}}{\delta}\right),$$

so that  $\rho_{\delta}$  also has integral 1, but vanishes outside the ball  $B(0, \delta)$ .

Now we consider the convolution product

$$u * \rho_{\delta}(\mathbf{x}) = \int_{\mathbb{R}^n} u(\mathbf{y}) \rho_{\delta}(\mathbf{x} - \mathbf{y}) d^n \mathbf{y}.$$

This is defined for all  $x \in \Omega$  with a distance less than  $\delta$  to the complement of  $\Omega$ . Thus, for any  $x \in \Omega$ , we can make  $\delta$  small enough so that  $u * \rho_{\delta}$  is defined at x.

Moreover,  $u * \rho_{\delta}$  is inifinetly differentiable: This is proved by differentiating with respect to the components of *x* under the integral sign, as much as you like.

Finally, the mean value property of u and the radial symmetry of  $\rho_{\delta}$  combine to ensure that  $u(\mathbf{x}) = u * \rho(\mathbf{x})$  for all  $\mathbf{x}$  where  $u * \rho$  is defined, which is what we were going to prove.

For a detailed argument, write

$$u * \rho(\mathbf{x}) = \int_{\mathbb{R}^n} u(\mathbf{x} - \mathbf{y}) \rho_{\delta}(\mathbf{y}) \, \mathrm{d}^n \mathbf{y}$$

and write the integral in polar form:

$$u * \rho(\mathbf{x}) = \int_0^\delta \int_{\partial B(\mathbf{x},r)} u(\mathbf{x} - \mathbf{y}) \rho_\delta(\mathbf{y}) \, \mathrm{d}S(\mathbf{y}) \, \mathrm{d}r$$
$$= \int_0^\delta r^{n-1} \int_{\mathbb{S}^{n-1}} u(\mathbf{x} - r\mathbf{y}) \rho_\delta(r\mathbf{y}) \, \mathrm{d}S(\mathbf{y}) \, \mathrm{d}r$$

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Now use the radial symmetry:  $\rho_{\delta}(ry)$  is constant for  $y \in \mathbb{S}^{n-1}$ , so this factor can be moved outside the inner integral. Next, use the mean value property of *u*:

$$\int_{\mathbb{S}^{n-1}} u(\boldsymbol{x} - r\boldsymbol{y})\rho_{\delta}(r\boldsymbol{y}) \,\mathrm{d}^n\boldsymbol{y} = A_n u(\boldsymbol{x}).$$

But u(x) is a constant, which we move outside the outer integral. We are left with

$$u * \rho(\mathbf{x}) = u(\mathbf{x}) \int_0^\delta r^{n-1} A_n \rho_\delta(r\mathbf{y}) \,\mathrm{d}r$$

where y is any unit vector. But running the whole calculation in reverse, this time without the *u* term, reveals that the integral here is merely the integral of  $\rho_{\delta}$ , which has the value 1. Thus we are left with  $u * \rho(x) = u(x)$ , as claimed.

## The maximum principle

**Definition.** A  $C^2$  function u is called *subharmonic* if  $\Delta u \ge 0$ , and *superharmonic* if  $\Delta u \le 0$ . Thus it is harmonic if and only if it is both subharmonic and superharmonic. (The reason for the naming will become clear later.) Obviously, u is superharmonic if and only if -u is subharmonic.

**Theorem 5** (Strong maximum principle). Assume that  $u \in C^2(\Omega)$  is subharmonic in a region  $\Omega \subseteq \mathbb{R}^n$ . If u has a global maximum in  $\Omega$ , then u is constant.

*Proof.* Let *M* be the global maximum of *u*, and put

$$S = \{ \boldsymbol{x} \in \Omega \mid u(\boldsymbol{x}) = M \}.$$

Then *S* is a closed subset of  $\Omega$ , by the continuity of *u*. It is also nonempty by assumption.

Consider any  $\mathbf{x} \in S$ . From (4) and the subharmonicity of u, we get  $\tilde{u}'_{\mathbf{x}}(r) \ge 0$  for r > 0. Thus we get  $\tilde{u}_{\mathbf{x}}(r) \ge \tilde{u}_{\mathbf{x}}(0) = u(\mathbf{x}) = M$  for r > 0 (so long as  $\overline{B}(\mathbf{x}, r) \subset \Omega$ ). But  $u \le M$  everywhere, and if u < M anywhere on the sphere  $\partial B(\mathbf{x}, r)$ , we would get  $\tilde{u}_{\mathbf{x}}(r) < 0$ . Thus u is contant equal to M on  $\partial B(\mathbf{x}, r)$  for any small enough r, and therefore, u is constant in some neighbourhood of  $\mathbf{x}$ . This means that S is open.

Since *S* is an open, closed, and nonempty subset of the connected set  $\Omega$ , we must have  $S = \Omega$ , and the proof is complete.

**Remark.** Obviously, we obtain a strong *minimum* principle for *superharmonic* functions by multiplying by -1. In particular, a non-constant harmonic function cannot attain a minimum or maximum value anywhere in  $\Omega$ .

**Corollary 6** (Weak maximum principle). Let  $\Omega \subset \mathbb{R}^n$  be a bounded domain, and assume that  $u \in C(\overline{\Omega}) \cap C^2(\Omega)$  is subharmonic. Then

$$\max\{u(\boldsymbol{x}) \mid \boldsymbol{x} \in \overline{\Omega}\} = \max\{u(\boldsymbol{x}) \mid \boldsymbol{x} \in \partial\Omega\}.$$

In particular, a harmonic function which is continuous on  $\overline{\Omega}$  attains its minimum and maximum values on the boundary  $\partial \Omega$ .

*Proof.* The weak principle is an obvious consequence of the strong principle. However, it is worth noting that a much more elementary proof exists. Namely, for any  $\varepsilon > 0$ , let  $v(\mathbf{x}) = u(\mathbf{x}) + \varepsilon |\mathbf{x}|^2$ , and note that then  $\Delta v > 0$ . But  $\Delta v(\mathbf{x}) \le 0$  if  $\mathbf{x}$  is an interior minimum point for v, so v cannot have any maximum in the interior. Thus for any  $\mathbf{x} \in \Omega$ ,

$$u(\mathbf{x}) = v(\mathbf{x}) - \varepsilon |\mathbf{x}|^2 \le \max_{\partial \Omega} v \le \max_{\partial \Omega} u + \varepsilon \max_{\mathbf{x} \in \partial \Omega} |\mathbf{x}|^2.$$

Now let  $\varepsilon \to 0$  to arrive at the conclusion  $u(\mathbf{x}) \leq \max_{\partial \Omega} u$ .

Our next result explains the terms *sub-* and *super*harmonic: A subharmonic function is below, and a superharmonic above, a harmonic function given the same boundary data.

**Corollary 7.** Assume that  $\Omega$  is a bounded domain, that  $u, v \in C(\overline{\Omega}) \cap C^2(\Omega)$ , that u is harmonic in  $\Omega$ , and that v = u on  $\partial \Omega$ . If v is subharmonic, then  $v \leq u$  in  $\Omega$ , while if v is superharmonic, then  $v \geq u$  in  $\Omega$ .

*Proof.* Apply the weak maximum principle to v - u if v is subharmonic, or to u - v if v is superharmonic.

**Remark.** Corollary 7 suggests a strategy for proving existence of a solution to the Dirichlet problem for the Laplace equation: Assume we are trying to find a harmonic function on  $\Omega$  with the given boundary value g. Consider the pointwise *supremum* of all subharmonic functions which are  $\leq g$  on  $\partial \Omega$ , and the pointwise *infimum* of all superharmonic functions which are  $\geq g$  on  $\partial \Omega$ . If the two functions coincide, they should provide a solution to the problem. This is the basis for *Perron's method*, which we will hopefully get a look at later.