

TMA4220 Numerical Solution of Partial Differential Equations Using Element Methods Fall 2019

Solutions to exercise set 1

1 We known that $\phi_i(\mathbf{x}_j) = \delta_{ij}$, where $\mathbf{x}_j, j = 1, 2, 3$ are the three corners, and we know that the linear basis functions are on the form

$$\phi_i(x,y) = a_i x + b_i y + c_i.$$

For ϕ_1 we get the system

$$\phi_i(1/2, 1/2) = 1/2a_1 + 1/2b_1 + c_1 = 1,$$

$$\phi_i(3/2, 1/2) = 3/2a_1 + 1/2b_1 + c_1 = 0,$$

$$\phi_i(1, 1) = a_1 + b_1 + c_1 = 0.$$

If we set up the same systems for ϕ_2 and ϕ_3 we get

[1/2]	1/2	1]	$\left\lceil a_{1}\right\rceil$	a_2	a_3		[1	0	0]
3/2	1/2	1	b_1	b_2	b_3	=	0	1	0
$\lfloor 1$	1	1	c_1	c_2	c_3		0	0	1

with solution

$$\phi_1(x, y) = -x - y + 2,$$

 $\phi_2(x, y) = x - y,$
 $\phi_3(x, y) = 2y - 1.$

2 We first start with the basis functions that are 1 in the corners, i.e. ϕ_1 , ϕ_2 and ϕ_3 . For these basis functions, we can write them on the form $\phi_i = d_i(\lambda_i - a_i)(\lambda_i - b_i)(\lambda_i - c_i)$ (why?). As we want the functions to be zero in all other nodes, we must have, for $i = 1, a_1 = 0, b_1 = 1/3, c_1 = 2/3$. Then, since $\phi_1(1,0,0) = 1, d_1 = (1-0)(1-1/3)(1-2/3) = 2/9$. Thus

$$\hat{\phi}_1 = 9/2\lambda_1(\lambda_1 - 1/3)(\lambda_1 - 2/3).$$

By symmetry, we find

$$\hat{\phi}_2 = 9/2\lambda_2(\lambda_2 - 1/3)(\lambda_2 - 2/3),$$

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For $\hat{\phi}_4$ we note that this has to be zero along the lines through $\hat{\phi}_1$ and $\hat{\phi}_2$, and through $\hat{\phi}_5$ and $\hat{\phi}_6$. Also, $\hat{\phi}_4(0, 2/3, 1/3) = 2/27$, so $d_4 = 27/2$. Thus, we have

$$\hat{\phi}_4 = 27/2\lambda_2\lambda_3(\lambda_3 - 1/3).$$

By symmetry,

$$\begin{split} \hat{\phi}_5 &= 27/2\lambda_2\lambda_3(\lambda_2 - 1/3), \\ \hat{\phi}_6 &= 27/2\lambda_3\lambda_1(\lambda_1 - 1/3), \\ \hat{\phi}_7 &= 27/2\lambda_3\lambda_1(\lambda_3 - 1/3), \\ \hat{\phi}_8 &= 27/2\lambda_1\lambda_2(\lambda_2 - 1/3), \\ \hat{\phi}_9 &= 27/2\lambda_1\lambda_2(\lambda_1 - 1/3). \end{split}$$

We see that $\hat{\phi}_{10}$ must be zero along all the edges, which yields

$$\hat{\phi}_{10} = 27\lambda_1\lambda_2\lambda_3.$$

a) By the fundamental theorem of calculus,

$$|u(x)|^{2} = \Big| \int_{-M}^{x} \frac{\partial u}{\partial x_{1}}(y_{1}, x_{2}) dy_{1} \Big|^{2},$$

and using Cauchy's inequality with $f = \frac{\partial u}{\partial x_1}(y_1, x_2)$ and g = 1, we get

$$\begin{split} \left| \int_{-M}^{x} \frac{\partial u}{\partial x_{1}}(y_{1}, x_{2}) dy_{1} \right|^{2} &\leq \int_{-M}^{x} \left| \frac{\partial u}{\partial x_{1}}(y_{1}, x_{2}) \right|^{2} dy_{1} \cdot \int_{-M}^{x} dx \\ &\leq 2M \int_{-M}^{M} \left| \frac{\partial u}{\partial x_{1}}(y_{1}, x_{2}) \right|^{2} dy_{1}. \end{split}$$

b) Integrate over x_2 to get

$$\int_{-M}^{M} |u(x)|^2 dx_2 \le 2M \int_{-M}^{M} \int_{-M}^{M} \left| \frac{\partial u}{\partial x_1}(y_1, x_2) \right|^2 dy_1 dx_2 = 2M \int_{\Omega} \left| \frac{\partial u}{\partial x_1}(y) \right|^2 dy_1 dx_2$$

where the equality comes from Fubini-Tonell's theorem.

c) Integrate over x_1 so

$$\begin{split} \|u\|_{L^{2}(\Omega)}^{2} &= \int_{\Omega} |u(x)|^{2} dx = \int_{-M}^{M} \int_{-M}^{M} |u(x)|^{2} dx_{2} dx_{1} \\ &\leq 2M \int_{-M}^{M} \int_{\Omega} \left| \frac{\partial u}{\partial x_{1}}(y) \right|^{2} dy dx_{1} \\ &= 4M^{2} \int_{\Omega} \left| \frac{\partial u}{\partial x_{1}}(y) \right|^{2} dy \\ &\leq 4M^{2} \int_{\Omega} \left| \frac{\partial u}{\partial x_{1}}(y) \right|^{2} + \left| \frac{\partial u}{\partial x_{2}}(y) \right|^{2} dy \\ &= 4M^{2} |u|_{H^{1}(\Omega)}^{2}. \end{split}$$

d) Note that since Ω is bounded its diameter is finite. Let $D = \operatorname{diam}(\Omega) < \infty$. Choose a point $x \in \Omega$ and let L = ||x||. Then $\Omega \subset [-(D+L), (D+L)]^2$. The function \bar{u} defined as

$$\bar{u}(x) = \begin{cases} u(x) & \text{if } x \in \Omega\\ 0 & \text{else} \end{cases}$$

satisfies the requirements.

Remarks:

- The Poincaré inequality on bounded domains $\Omega \subset \mathbb{R}^d$ is shown in an analogous fashion.
- Note that the constant C grows as the box becomes larger.
- If $u \in H_0^k(\Omega)$ for any $k \ge 1$, then $\frac{\partial u}{\partial x_i} \in H_0^{k-1}(\Omega)$ for any $i = 1, \ldots, d$, and more generally, $D^{\alpha}u \in H_0^{k-|\alpha|}$ for any multiindex α of size $\alpha \le k$. Thus, we can iterate Poincaré inequality and find that there is a constant C > 0 such that

$$||u||_{L^2(\Omega)} \le C|u|_{H^1(\Omega)} \le \dots \le C|u|_{H^k(\Omega)}, \forall u \in H_0^k(\Omega).$$