# Project I – Helmholtz equation in two dimensions

The deadline for this assignment is 28. October 2018 and counts towards 15% of the final grade. The delivery consists of a report answering all the questions and presenting the results, together with a source code in MATLAB. Working in pairs is possible but not compulsory.

#### 1 Problem formulation

Let us consider the Helmholtz equation posed on  $\Omega = (0,1)^2$ ,

$$u(\mathbf{x}) - \Delta u(\mathbf{x}) = f(\mathbf{x}), \quad \forall \mathbf{x} = (x, y) \in \Omega$$
 (1)

with source term  $f(x) = (2\pi^2 + 1)\cos(\pi x)\sin(\pi y)$ , and supplemented with boundary conditions,

$$u(\boldsymbol{x}) = u_D , \quad \boldsymbol{x} \in \Gamma_D = \{ \boldsymbol{x} \in \partial\Omega : \boldsymbol{y} \in \{0, 1\} \}$$
  
$$\nabla u \cdot \mathbf{n}(\boldsymbol{x}) = 0 , \quad \boldsymbol{x} \in \Gamma_N = \{ \boldsymbol{x} \in \partial\Omega : \boldsymbol{x} \in \{0, 1\} \}$$
 (2)

with **n** the outward normal to the boundary  $\partial\Omega$ .

The goal of this assignment is to write an algorithm to compute approximate solutions to Problem (1)–(2) using linear Lagrange Finite Elements on a triangular mesh  $\mathcal{T}_h$ .

- (a) Verify that  $\tilde{u}:(x,y)\mapsto\cos(\pi x)\sin(\pi y)$  is solution to (1)–(2).
- (b) Derive a weak formulation of (1)–(2), specify the function spaces.
- (c) Is the solution  $\tilde{u}$  unique?

#### 2 Finite Element space

The construction of the approximation space based on Lagrange  $\mathbb{P}_1$  Finite Elements is now considered.

(a) Give the definition of the Lagrange  $\mathbb{P}_1$  reference element on the unit triangle  $\hat{K}$  with vertices  $\{\hat{\mathbf{v}}_0 = (0,0), \hat{\mathbf{v}}_1 = (1,0), \hat{\mathbf{v}}_2 = (0,1)\}$ , and associated local shape functions  $(\hat{\varphi}_0, \hat{\varphi}_1, \hat{\varphi}_2)$ , then implement it

- (b) Write a simple test showing that shape functions  $(\hat{\varphi}_0, \hat{\varphi}_1, \hat{\varphi}_2)$  form a nodal basis, and that for any  $\hat{x} \in \hat{K}$ ,  $\hat{\varphi}_0(\hat{x}) + \hat{\varphi}_1(\hat{x}) + \hat{\varphi}_2(\hat{x}) = 1$ .
- (c) Implement the affine mapping  $T_K : \hat{K} \to K$  and verify that the determinant of the Jacobian  $J_{T_K}$  is positive for triangle  $K_0 = \{(1,0), (3,1), (3,2)\}$ . Interpret this result.
- (d) Implement the inverse of  $J_{T_K}$  and verify that the Finite Element obtained by transporting the reference Finite Element  $(\hat{K}, \hat{\mathcal{P}}, \hat{\Sigma})$  to  $K_0$  is equivalent to  $(\hat{K}, \hat{\mathcal{P}}, \hat{\Sigma})$ .
- (e) Formulate the Galerkin problem corresponding to the weak formulation derived at the preceding section, in particular define the approximation space carefully.

#### 3 Numerical integration

Unless efficient exact evaluation is possible, computation of integrals is performed using quadrature rules. Such approximations are expressed as the weighted sum of integrand values over  $N_q$  quadrature points  $\{\zeta_q\}$ ,

$$I_K = \int_K \psi(\boldsymbol{x}) \, \mathrm{d}\boldsymbol{x} \approx \sum_{q=1}^{N_q} \psi(\boldsymbol{\zeta}_q) \; \omega_q$$

with real coefficients  $\{\omega_q\}$  called quadrature weights. The order  $k_q$  of a quadrature rule is the polynomial degree of the integrand for which the evaluation is exact. In particular, Gauss-Legendre quadratures on a real interval gathered in Table 3 satisfy the relation  $k_q = 2N_q - 1$ . Quadrature rules can be defined using other polynomials and considering higher dimensions in space. In the frame of Finite Elements, contributions on the reference simplex  $\hat{K}$  can be written as

$$\int_{\hat{K}} \hat{\psi}(\hat{\boldsymbol{x}}) \, \mathrm{d}\hat{\boldsymbol{x}} \approx \sum_{q=1}^{N_q} \hat{\psi}(\boldsymbol{\zeta}_q) \, \hat{\omega}_q$$

with q the index of the quadrature point. Therefore any contribution on cell  $K \in \mathcal{T}_h$  is obtained directly by composition with the affine change of coordinates  $T_K$ ,

$$\int_{K} \psi(\boldsymbol{x}) \, \mathrm{d}\boldsymbol{x} \approx |\det(J_{T_{K}})| \sum_{q=1}^{N_{q}} \psi \circ T_{K}(\boldsymbol{\zeta}_{q}) \, \hat{\omega}_{q}$$

with  $J_{T_K}$  elementwise constant. If  $\psi$  involves derivatives, the change of variable should take into account that  $(f \circ g)' = (f' \circ g) \cdot g'$ .

(a) Implement Gauss-Legendre quadratures from Table 3 and plot the approximation error for

$$I = \int_{1}^{2} e^{x} \, \mathrm{d}x$$

$k_q$	$N_q$	$\{\hat{oldsymbol{\zeta}}_q\}$	$\{\hat{\omega}_q\}$
1	1	$ar{\zeta}$	I
3	2	$\bar{\zeta} \pm  I  \frac{\sqrt{3}}{6}$	$\frac{1}{2} I $
5	3	$\bar{\zeta} \pm  I  \frac{\sqrt{15}}{10}$	$rac{5}{18} I $
		$ar{\zeta}$	$\frac{8}{18} I $
7	4	$\bar{\zeta} \pm  I  \frac{\sqrt{525 + 70\sqrt{30}}}{70}$	$\frac{18-\sqrt{30}}{36} I $
		$\bar{\zeta} \pm  I  \frac{\sqrt{525 - 70\sqrt{30}}}{70}$	$\frac{18+\sqrt{30}}{36} I $

Table 3: Gauss–Legendre quadratures on the interval [a,b] with  $\bar{\zeta}=(a+b)/2,$  and |I|=|b-a|

$k_q$	$N_q$	$\{\hat{oldsymbol{\zeta}}_q\}$	$\{\hat{\omega}_q\}$
1	1	$\left(\frac{1}{3},\frac{1}{3},\frac{1}{3}\right)$	K
2	3	$\left(\frac{1}{2},\frac{1}{2},0\right)_3$	$\frac{1}{3} K $
3	4	$\left(\frac{1}{3},\frac{1}{3},\frac{1}{3}\right)$	$\frac{-9}{16} K $
		$\left(\frac{1}{5},\frac{1}{5},\frac{3}{5}\right)_3$	$\frac{25}{48} K $
4	7	$\left(\frac{1}{3},\frac{1}{3},\frac{1}{3}\right)$	$\frac{9}{40} K $
		$(a_i, a_i, 1 - 2a_i)_3$	$\frac{155\pm\sqrt{15}}{1200} K $
		$a_i = \frac{6 \pm \sqrt{15}}{21}$	

Table 3: Gauss–Legendre quadratures on a triangle K in barycentric coordinates  $\{\hat{\zeta}_q\} = (\lambda_0, \lambda_1, \lambda_2)$ , with  $(\cdot, \cdot, \cdot)_k$  the k distinct tuples obtained by permutation, [1] page 360.

(b) Implement Gauss-Legendre quadratures from Table 3 and plot the approximation error for

$$I = \int_{K_0} \log(x+y) \, \mathrm{d}x$$

with  $K_0 = \{(1,0), (3,1), (3,2)\}.$ 

(c) Discuss why the choice of quadrature is important for the evaluation of Finite Element contributions. Which properties of the problem should be considered for terms corresponding to the left-hand side and right-hand side of the equation?

#### 4 Assembly of the linear system

For each cell  $K \in \mathcal{T}_h$ , elementwise contributions for the Helmholtz equation are under the form of a sum of two submatrices, corresponding to contributions of the mass matrix

$$\mathbf{M}_K = \left[ \int_K \varphi_j(\boldsymbol{x}) \varphi_i(\boldsymbol{x}) \, \mathrm{d} \boldsymbol{x} \right]_{ij}$$

and of the stiffness matrix

$$\mathbf{K}_K = \left[ \int_K \nabla \varphi_j(\hat{\boldsymbol{x}}) \nabla \varphi_i(\boldsymbol{x}) \, \mathrm{d} \boldsymbol{x} \right]_{ij}$$

with j indices of the global shape functions (solution space), and i indices of the global basis functions (test space) with support on K. For any  $K \in \mathcal{T}_h$  assembling the local equation consists of computing the contributions for indices  $\hat{j} = 1, \dots, N_{\mathcal{P}}$  of the local shape functions (solution space), and  $\hat{i} = 1, \dots, N_{\mathcal{P}}$  indices of the local basis functions (test space) with support on K,  $N_{\mathcal{P}}$  the dimension of the Finite Element. The passage from one to another is performed with a mapping from (cell) local indices  $(\hat{i}, \hat{j})$  to (mesh) global indices (i, j). The obtained submatrix and subvector are then added to the global matrix and load vector.

- (a) Detail the assembly of the local matrix and the local vector for any  $K \in \mathcal{T}_h$ .
- (b) Describe the assembly of the Dirichlet and Neumann boundary conditions.

# 5 Convergence analysis

(a) Implement the computation of the  $L^2$  error norm given by

$$\|u - u_h\|_{\mathrm{L}^2} = \left(\int_{\Omega} |u(\boldsymbol{x}) - u_h(\boldsymbol{x})|^2 d\boldsymbol{x}\right)^{\frac{1}{2}}$$

Why should you be careful with the evaluation of the integral?

(b) Solve the problem for different mesh sizes  $h_{\mathcal{T}} = 1/M$  with M = 4, 8, 16 and plot the L<sup>2</sup> error norm with respect to the dimension of the problem.

### 6 Extension to an evolution problem

Let us consider the evolution problem,

$$\partial_t u(\boldsymbol{x},t) - \nu \Delta u(\boldsymbol{x},t) = f(\boldsymbol{x},t), \quad \forall (\boldsymbol{x},t) \in \Omega \times (0,T)$$
 (3)

with  $u(\boldsymbol{x},0)=u_0$  given initial data, and  $\nu$  diffusivity.

(a) Describe how you would modify the algorithm developed for the Helmholtz problem to solve this equation for a given discretization in time. For example use the Backward Euler scheme. The function  $\tilde{u}(\boldsymbol{x},t) = e^{-\nu t} \sin(x\cos(\theta) + y\sin(\theta))$  can be used to verify the implementation for the homogeneous equation (optional); take  $\nu = 1$  and  $\theta = \pi/4$  for instance.

## **Bibliography**

[1] A. Ern and J.-L. Guermond. Theory and Practice of Finite Elements, volume 159 of Springer Series: Applied Mathematical Sciences. Springer, 2004.