

Problem 7:

a) Scheme:

$$\frac{1}{\kappa} (U_i^{n+1} - U_i^n) = \frac{\alpha}{h^2} (U_{i+1}^{n+1} - 2U_i^{n+1} + U_{i-1}^{n+1}) - k \cdot U_i^{n+1},$$

=>

$$U_i^{n+1} = U_i^n + \alpha \frac{k}{h^2} (U_{i+1}^{n+1} - 2U_i^{n+1} + U_{i-1}^{n+1}) - k \cdot U_i^{n+1}, \quad i=1, \dots, M-1$$

$$U_0^n = U_M^n = 0 \quad n=0, 1, 2, \dots$$

Local truncation error, \tilde{e}_i^n : Let $u_i^n = u(x_i, t_n)$

$$k \cdot \tilde{e}_i^n = U_i^{n+1} - U_i^n - \alpha \frac{k}{h^2} (U_{i+1}^{n+1} - 2U_i^{n+1} + U_{i-1}^{n+1}) - k \cdot U_i^{n+1}$$

For simplicity, do the expansion around (x_i, t_{n+1}) :

$$k \cdot \tilde{e}_i^n = k \cdot \partial_t u + \frac{1}{2} k^2 \partial_t^2 u + O(k^3)$$

$$- \alpha k [\partial_x^2 u + h^2 \cdot \frac{1}{12} \partial_x^4 u + O(h^6)] + k \cdot u$$

so, using that $\partial_t u - \partial_x^2 u + u = 0$ we get

$$\tilde{e}_i^n = \frac{1}{2} k \partial_t^2 u - h^2 \frac{\alpha}{12} \partial_x^4 u + O(k^2) + O(h^4)$$

all evaluated in (x_i, t_{n+1}) , or in (x_i, t_n) .

The difference will only appear in the higher order terms.

b) The scheme is consistent ($\tilde{e}_i^n \rightarrow 0$), so according to the Lax equivalence theorem, it is sufficient to prove stability.

Let $r = \frac{k}{h^2}$. The scheme can be written as

$$U_i^{n+1} - r(U_i^{n+1} - 2U_i^n + U_{i-1}^{n+1}) + k \cdot U_i^{n+1} = U_i^n.$$

Using $U^n = [U_1^n, \dots, U_M^n]^T$ this becomes

$$A \cdot U^{n+1} = U^n \Rightarrow U^{n+1} = C \cdot U^n \text{ with } C = A^{-1}, \text{ and}$$

where $A = (1+k)I + r \cdot T$, $T = \text{tridiag}\{-1, 2, -1\}$

We know: The scheme is stable if $\rho(C) \leq 1$

The eigenvalues are:

$$\lambda_j(T) = 4 \sin^2\left(\frac{j\pi}{2M}\right) \quad (\text{from the note})$$

$$\lambda_j(A) = (1+k) + r \cdot \lambda_j(T) = 1+k + 4r \sin^2\left(\frac{j\pi}{2n}\right)$$

and

$$\lambda_j(c) = \frac{1}{\lambda_j(A)} = \frac{1}{1+k + 4r \sin^2\left(\frac{j\pi}{2n}\right)}$$

and clearly

$$0 < \lambda_j(c) < 1 \Rightarrow \rho(c) < 1$$

and the scheme is stable, and thus convergent.

Problem 2:

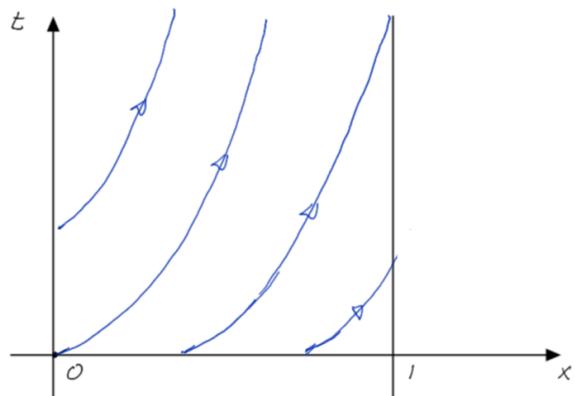
a) The characteristics are the solution of the equations

$$\frac{dx}{dt} = \frac{1}{(1+t)^2} \quad x(t^*) = x^*.$$

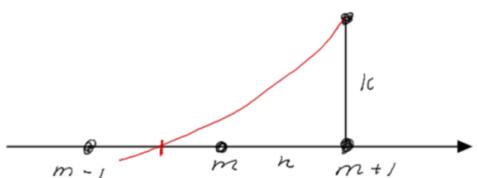
The solution is

$$x(t) = x^* + \frac{1}{1+t^*} - \frac{1}{1+t}$$

The characteristics are moving from left to right, so we should include boundary conditions at $x=0$.



b) We notice that the velocity $\frac{1}{(1+t)^2}$ becomes smaller as t increases, so for the CFL condition to be satisfied everywhere, it suffices to have it satisfied for the first step. The requirement then is that the characteristic through (x_{m+1}, t_0) should pass through the interval (x_{m-1}, x_{m+1}) at $t=0$ (see the figure).



Computational stencil

The condition becomes:

$$\begin{aligned} x(0) &= x_{m+1} + \frac{1}{1+t_0} - 1 \in [x_{m-1}, x_{m+1}] \\ \Rightarrow x_{m+1} + \frac{1}{1+t_0} - 1 &\geq x_{m-1} \\ \Rightarrow 2h + \frac{1}{1+t_0} &\geq 1 \Rightarrow t_0 \leq \frac{2h}{1-2h} \end{aligned}$$

If $h = 1/4$, $t_0 \leq 1$.

Problem 3:

This is not a complete solution to this problem, but some indications on what can be done:

If we do not know much about ν and a , it may be a good idea to go for 2. order approximations:

$$\begin{aligned} \nu(u) u_x &\rightarrow \nu(u_i) \cdot \frac{1}{2h} (U_{i+1} - U_{i-1}) \\ (a(u) u_x)_x &\rightarrow h^{\frac{1}{2}} (a_{i+\frac{1}{2}} U_{i+1} - (a_{i+\frac{1}{2}} + a_{i-\frac{1}{2}}) U_i + a_{i-\frac{1}{2}} U_{i-1}) \end{aligned}$$

where one of the following can be used:

$$a_{i+\frac{1}{2}} = \frac{1}{2} (a(U_{i+1}) + a(U_i))$$

or

$$a_{i+\frac{1}{2}} = a\left(\frac{1}{2}(U_{i+1} + U_i)\right)$$

Boundary conditions:

$$x = 0 : U_0 = 0$$

$$\begin{aligned} x = 1 : u_x + u = 0 &\rightarrow \frac{U_{M+1} - U_{M-1}}{2h} + U_M = 0 \\ \Rightarrow U_{M+1} &= U_{M-1} - 2h U_M \end{aligned}$$

which is inserted into the difference formula (not written out here) for $i = M$.

Altogether, this ends up in a system of ODEs.

$$\ddot{U} = F(U), \text{ where } U = [U_1, \dots, U_M].$$

Since we have a diffusion term, we would choose an implicit (A-stable) scheme to solve the ODE. The trapezoidal rule (Crank-Nicolson) is of order 2, so that may be a good choice:

$$U^{n+1} = U^n + \frac{h}{2} [F(U^{n+1}) + F(U^n)].$$

which altogether creates a system of nonlinear equations which has to be solved for each step.

Other considerations:

If the diffusion term is small, you may want to use an upwind scheme ($\nu > 0$):

$$\nu(u) u_x \rightarrow \nu(u_i) \frac{U_i - U_{i-1}}{h}. \quad (\text{order 1})$$

In this case, it may also be possible to use an explicit scheme, but stability may be an issue.

Problem 4: $\alpha > 0$

$$\text{Scheme: } U_m^{n+1} = U_m^n - \rho \cdot (U_m^n - U_{m-1}^n)$$

Let $U_m^n = g^n e^{i\beta x_m}$, and we get

$$g = 1 - \rho(1 - e^{-i\beta h}) = 1 - \rho + \rho(\cos(\beta h) - i \sin(\beta h))$$

and

$$\begin{aligned} |g|^2 &= (1 - \rho + \rho \cos(\beta h))^2 + \rho^2 \sin^2(\beta h) \\ &= 1 - 2\rho(1-\rho)(1-\cos\beta h) \geq 0, \quad -\pi < \beta h \leq \pi \end{aligned}$$

So, for $\rho = 1$, $|g|^2 = 1$ and there are no dissipation

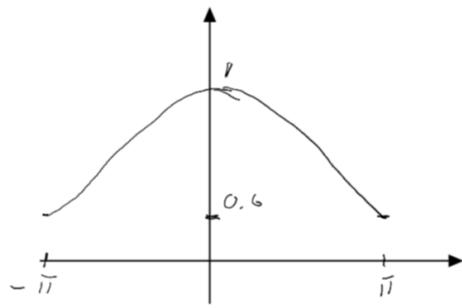
$\rho > 1$ or $\rho < 0$: $|g|^2 > 0$ for some βh , and the scheme is unstable. This coincide with the CFL-condition.

$0 < \rho < 1$: In this case, $|g| \leq 1$ and the scheme is stable. g depends on βh , so the method is dissipative. It can be shown that

$$|g| = 1 - \frac{1}{2}\rho(\rho-1)(\beta h)^2 + O((\beta h)^4)$$

thus the method is dissipative of order 2.

For $\rho = 0.8$, $g = \sqrt{0.68 + 0.3 \cdot \cos(\beta h)}$, is given by

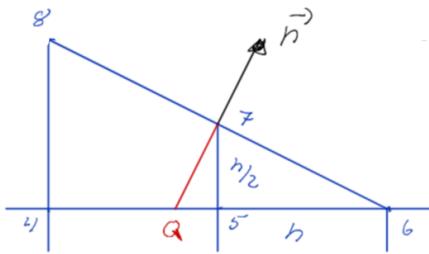


and high frequency modes will be damped faster than slow modes. Consequently, the solution will be smoothed out over time.



Problem 5 :

An approximation to the solution in point Q can be found by linear interpolation:



The distance from Q to node 5 is

$$h' = \frac{h}{4}.$$

$$\begin{aligned} u_Q &\approx u_4 + \frac{u_5 - u_4}{h} \cdot \frac{3}{4} h \\ &= \frac{3}{4} u_5 + \frac{1}{4} u_4 \end{aligned}$$

The distance from node F to Q is $\frac{\sqrt{5}}{4}h$, so a suitable approximation to $\partial_n u$ in (7) may be

$$\frac{u_7 - u_Q}{\frac{\sqrt{5}}{4}h} = g_F \Rightarrow \frac{4}{\sqrt{5}h} (u_7 - \frac{3}{4}u_5 - \frac{1}{4}u_4) = g_F$$

$$\Rightarrow \boxed{\frac{1}{\sqrt{5}h} (4u_7 - 3u_5 - u_4) = g_F}$$

To find the local truncation error, notice that the unit normal vector $n\text{-hat}$ is

$$n\text{-hat} = \left[\frac{1}{\sqrt{5}}, \frac{2}{\sqrt{5}} \right]^T \quad \text{so}$$

$$\partial_n u = n\text{-hat} \cdot \nabla u = \frac{1}{\sqrt{5}}(u_x + 2u_y).$$

Then

$$\begin{aligned} \varepsilon_F &= \frac{1}{\sqrt{5}h} (4u(x_F, y_F) - 3u(x_F, y_F - \frac{h}{2}) - u(x_F - h, y_F - \frac{h}{2})) - g_F \\ &= \frac{1}{\sqrt{5}h} (4u - 3(u - \frac{h}{2}u_y + \frac{1}{2}(\frac{h}{2})^2 u_{yy}) + O(h^3)) \\ &\quad - (u - hu_x - \frac{h}{2}u_y + \frac{1}{2}h^2 u_{xx} + \frac{1}{2}h^2 u_{xy} + \frac{1}{2}(\frac{h}{2})^2 u_{yy} + O(h^3)) \\ &= \underbrace{\frac{1}{\sqrt{5}}(u_x + 2u_y)}_{=0} - g_F - \frac{h}{2}(u_{xx} + u_{yy} + u_{xy}) + O(h^2) \end{aligned}$$

so the local truncation error is of order 1.

Problem 6.

$$-(a(x) \cdot u_x)_x = 1 \quad 0 < x < 1 \\ u_x(0) = 1, \quad u(1) = 0.$$

a) Multiply by a function v , use partial integration to get

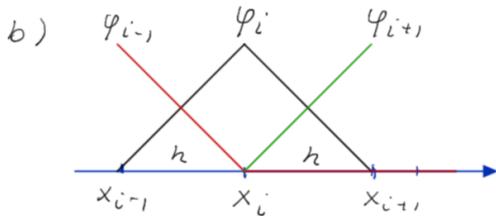
$$\int_0^1 a u_x v_x dx - \int_0^1 u_x v = \int_0^1 v dx$$

for the integrals to exist, $u, v \in H^1(\Omega)$, $\Omega = (0, 1)$. We know that $u_x(0) = 1$. To get rid of the boundary term at $x=1$, we require $v(1) = 0$ (and we know that $u(1) = 0$).

$$\text{let } V = H_n^1(\Omega) = \{v \in H^1(\Omega), v(1) = 0\} \quad (\Omega = (0, 1)),$$

The weak formulation is now:

Find $u \in H_n^1$ s.t. $\int_0^1 a u_x v_x dx = \int_0^1 v dx - v(1), \forall v \in H_n^1(\Omega)$



The linear nodal basis functions are described in the figure.

$$\varphi_i' = \begin{cases} 1/h & x_{i-1} < x < x_i \\ -1/h & x_i \leq x < x_{i+1} \\ 0 & \text{otherwise.} \end{cases}$$

The elements of the stiffness matrix is

$$A_{ij} = \int_0^1 (x + \frac{1}{2}) \varphi_i' \varphi_j' dx, \quad i = 0, \dots, M-1$$

The matrix is tridiagonal, with elements: $i=1, \dots, M-1$

$$A_{ii} = \frac{1}{h^2} \int_{x_{i-1}}^{x_{i+1}} (x + \frac{1}{2}) dx = \frac{1}{h} (2x_i + 1)$$

$$A_{i,i+1} = -\frac{1}{h^2} \int_{x_i}^{x_{i+1}} (x + \frac{1}{2}) dx = -\frac{1}{h} (x_i + \frac{1}{2}(1+h))$$

$$A_{i-1,i} = -\frac{1}{h} (x_i + \frac{1}{2}(1-h))$$

For the boundary function φ_0 , we get.

$$A_{00} = \frac{1}{h^2} \int_0^h (x + \frac{1}{2}) dx = \frac{1}{2h} (1+h)$$

$$A_{01} = -\frac{1}{h^2} \int_0^h (x + \frac{1}{2}) dx = -\frac{1}{2h} (1+h)$$

For the load vector \vec{b} , all elements
but the first is

$$b_i = \int_{x_{i-1}}^{x_{i+1}} l \cdot \varphi_i(x) dx = 2 \cdot \frac{1}{2} h = h$$

For the first one we get

$$b_0 = \int_0^h l \cdot \varphi_0(x) dx + \varphi_0(1) = \frac{h}{2} + 1$$