



1 **B&N: 3.7** Due to figures, the exercise solution is found at the end of this solution proposal.

2 **B&N: 3.14**

- a) If we consider  $n$ -periodic sequences  $u_k$ , the sequence  $L[u]_k$  is also  $n$ -periodic as can be seen from the following:

$$L[u]_{k+n} = u_{k+n+1} - 2u_{k+n} + u_{k+n-1} = u_{k+1} - 2u_k + u_{k-1} = L[u]_k.$$

- b) Since  $u_k$  is 4-periodic, it is represented by the vector  $[u_1 \ u_2 \ u_3 \ u_4]^T$ . The same goes for  $L[u]_k$ , which is characterized by the vector  $[L[u]_1 \ L[u]_2 \ L[u]_3 \ L[u]_4]^T$ . Furthermore, we can see that

$$[L[u]_1 = u_2 - 2u_1 + u_0 = u_2 - 2u_1 + u_4$$

$$[L[u]_2 = u_3 - 2u_2 + u_1$$

$$[L[u]_3 = u_4 - 2u_3 + u_2$$

$$[L[u]_4 = u_5 - 2u_4 + u_3 = u_1 - 2u_4 + u_3.$$

In total, this means that

$$\begin{bmatrix} [L[u]_1 \\ [L[u]_2 \\ [L[u]_3 \\ [L[u]_4 \end{bmatrix} = \begin{bmatrix} -2 & 1 & 0 & 1 \\ 1 & -2 & 1 & 0 \\ 0 & 1 & -2 & 1 \\ 1 & 0 & 1 & -2 \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{bmatrix},$$

meaning  $L$  is represented by the matrix, called  $M_4$ .

- c) *This answer is quite long since we want to use the same procedure in part d).* We look for eigenvectors and  $\lambda$ -values of the matrix  $M_4$ . Since  $L$  is the linear transformation represented by  $M_4$ , this is the same as finding  $\lambda \in \mathbb{R}$  and  $u \in \mathbb{R}^4$  such that

$$L[u]_k = u_{k+1} - 2u_k + u_{k-1} = \lambda u_k, \quad k = 1, 2, 3, 4 \quad (1)$$

where  $u_k$  can be considered a sequence in  $S_4$ . We take the DFT of this to find

$$\begin{aligned} e^{\frac{2\pi i k}{4}} \hat{u}_k - 2\hat{u}_k + e^{-\frac{2\pi i k}{4}} \hat{u}_k &= \lambda \hat{u}_k \\ \Rightarrow \lambda_k &= 2 \left( \cos\left(\frac{2\pi k}{4}\right) - 1 \right), \quad k = 1, 2, 3, 4. \end{aligned}$$

That is, the eigenvalues are  $-2, -4, -2$ , and  $0$ . Relabeling the eigenvalues as  $\lambda_j$ , we reinsert the now known eigenvalues into (1) to find the components  $u_k^j$  of the  $j$ 'th

eigenvector:

$$\begin{aligned} u_{k+1}^j - 2u_k^j + u_{k-1}^j &= \lambda_j u_k^j \\ &= 2 \cos\left(\frac{2\pi j}{4}\right) u_k^j - 2u_k^j. \end{aligned}$$

That is, we need

$$u_{k+1}^j + u_{k-1}^j = 2 \cos\left(\frac{2\pi j}{4}\right) u_k^j = e^{i\frac{2\pi j}{4}} u_k^j + e^{-i\frac{2\pi j}{4}} u_k^j.$$

This should be enough information for us to make the ansatz that

$$u_k^j = C_j e^{i\frac{2\pi jk}{4}} = C_j w^{jk} \quad (w = e^{i\frac{2\pi}{4}}).$$

The constants  $C_j$  are arbitrary; recall that if  $v$  is an eigenvector of a matrix  $A$ , then  $Cv$  is also an eigenvector of  $A$  for  $C \in \mathbb{C}$ . One can easily check (e.g. by insertion into (1)) that the eigenvectors are, indeed, correct. Thus, we have the eigenvector-eigenvalue pairs  $(\lambda_j, u^j)$ , where

$$\lambda_j = 2 \left( \cos\left(\frac{2\pi j}{4}\right) - 1 \right), \quad u^j = \begin{bmatrix} u_1^j \\ u_2^j \\ u_3^j \\ u_4^j \end{bmatrix} = C_j \begin{bmatrix} w^j \\ w^{2j} \\ w^{3j} \\ w^{4j} \end{bmatrix}.$$

Now, if we choose the particular eigenvector representatives with  $C_j = w^{-j}$ , we have

$$u^1 = \begin{bmatrix} 1 \\ w \\ w^2 \\ w^3 \end{bmatrix} \quad u^2 = \begin{bmatrix} 1 \\ w^2 \\ w^4 \\ w^6 \end{bmatrix} \quad u^3 = \begin{bmatrix} 1 \\ w^3 \\ w^6 \\ w^9 \end{bmatrix} \quad u^4 = \begin{bmatrix} 1 \\ w^4 \\ w^8 \\ w^{12} \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \\ 1 \\ 1 \end{bmatrix}.$$

So, we can see that the eigenvectors are the column vectors of the DFT matrix  $F_4$  as seen in equation (3.2) in B&N! That is, we can rearrange the eigenvectors and eigenvalues to arrive at the eigenvalue decomposition  $M = F_4 \Lambda_4 F_4^{-1}$ , where

$$F_4 = \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & w & w^2 & w^3 \\ 1 & w^2 & w^4 & w^6 \\ 1 & w^3 & w^6 & w^9 \end{bmatrix}, \quad \Lambda_4 = \begin{bmatrix} \lambda_4 & 0 & 0 & 0 \\ 0 & \lambda_1 & 0 & 0 \\ 0 & 0 & \lambda_2 & 0 \\ 0 & 0 & 0 & \lambda_3 \end{bmatrix}.$$

This means that we can use the FFT to diagonalize  $M_4$ , as we shall now see. We write the diagonalization of  $M_4$  in terms of multiplication with  $F_4$  and its inverse:

$$M_4 = F_4 \Lambda_4 F_4^{-1} \Leftrightarrow F_4^{-1} M_4 = \Lambda_4 F_4^{-1} \Leftrightarrow (F_4^{-1} M_4)^H = (F_4^{-1})^H \Lambda_4 \Leftrightarrow ((F_4^{-1})^H)^{-1} (F_4^{-1} M_4)^H = \Lambda_4.$$

Here,  $A^H$  denotes the conjugate transpose of the matrix  $A$ . Since  $F_4$  is symmetric, conjugate transposition reduces to just conjugation, i.e.  $F_4^H = \overline{F_4}$ . Also, as a result from basic linear algebra, conjugate transposition commutes with inversion of matrices, i.e.  $(A^H)^{-1} = (A^{-1})^H$ . Using these two facts simplifies the diagonalization such that we have

$$\Lambda_4 = ((F_4^{-1})^H)^{-1} (F_4^{-1} M_4)^H = ((F_4^H)^{-1})^{-1} (F_4^{-1} M_4)^H = \overline{F_4} (F_4^{-1} M_4)^H.$$

Furthermore, from the remark after Theorem 3.3, we know that  $F_4^{-1} = \frac{1}{4}\overline{F_4}$ , such that

$$\Lambda_4 = \frac{1}{4}\overline{F_4}(F_4 M_4)^H. \quad (2)$$

Now, consider  $M_4$  expressed in terms of its column vectors  $m_1, m_2, m_3, m_4$ :

$$M_4 = [m_1|m_2|m_3|m_4].$$

Multiplying this by any matrix  $A$  is equivalent to multiplying its column vectors by  $A$ :

$$AM_4 = [Am_1|Am_2|Am_3|Am_4].$$

In particular, multiplying  $m_i$  with  $\overline{F_4}$  is equivalent to taking the DFT of  $m_i$ . Hence, we have

$$\overline{F_4}M_4 = [\overline{F_4}m_1|\overline{F_4}m_2|\overline{F_4}m_3|\overline{F_4}m_4] = [\hat{m}_1|\hat{m}_2|\hat{m}_3|\hat{m}_4].$$

The DFT can be computed using the efficient FFT algorithm. So, in total, we can compute the diagonalization as given by (2) in three steps:

- 1:** Compute the FFT of the columns of  $M_4$ .
- 2:** Transpose the resulting matrix.
- 3:** Compute the FFT of the columns of the transposed matrix and multiply by  $1/4$ .

*Remark:* This diagonalization procedure is powerful! A variant of it is used to compute finite difference approximations to the 2D Poisson equation with just  $\mathcal{O}(\log n)$  operations per unknown as shown in the note at

[https://www.math.ntnu.no/emner/TMA4205/2015h/notes/Poisson2D\\_diag.pdf](https://www.math.ntnu.no/emner/TMA4205/2015h/notes/Poisson2D_diag.pdf)

- d)** Generalizing the result from c) for all  $n$  is not too challenging; you need to follow the same steps, and find that

$$\lambda_k = 2 \left( \cos \left( \frac{2\pi k}{n} \right) - 1 \right), \quad k = 1, 2, \dots, n.$$

Then, by the same argument as earlier, we can (by ansatz) find that the components of the eigenvectors are

$$u_k^j = C_j e^{i \frac{2\pi jk}{n}} = C_j w^{jk} \quad (w = e^{i \frac{2\pi}{n}}).$$

Normalizing the eigenvectors with  $C_j = w^{-j}$  gives the eigenvalue decomposition of  $M_n$  as

$$M_n = S_n \Lambda_n S_n^{-1},$$

and the discussion follows the same track to end up with the diagonalization procedure:

- 1:** Compute the FFT of the columns of  $M_n$ .
- 2:** Transpose the resulting matrix.
- 3:** Compute the FFT of the columns of the transposed matrix and multiply by  $1/n$ .

3 **B&N: 3.16** We want to find the Z-transform of the sequence  $x = \{x_j\}_{j=-\infty}^{\infty}$  where

$$x_j = \begin{cases} \frac{1}{2^j}, & j \geq 0 \\ 0, & j < 0. \end{cases}$$

Using the definition of the Z-transform, we have

$$\hat{x}(\phi) = \sum_{j=-\infty}^{\infty} x_j e^{-ij\phi} = \sum_{j=0}^{\infty} \left(\frac{e^{-i\phi}}{2}\right)^j = \frac{1}{1 - \frac{e^{-i\phi}}{2}} = \frac{1 - \frac{e^{i\phi}}{2}}{\frac{5}{4} - \cos(\phi)}.$$

Alternatively, writing  $z = e^{i\phi}$ :

$$\hat{x}(z) = \sum_{j=-\infty}^{\infty} x_j z^{-j} = \sum_{j=0}^{\infty} \left(\frac{z}{2}\right)^j = \frac{1}{1 - \frac{z}{2}} = \frac{1 - \bar{z}}{\frac{5}{4} - \operatorname{Re}(z)}.$$

4 **B&N: 3.7** The figures below show  $f$  and the filtered versions of  $f$  with varying values of  $m$ . Two observations can be made. The first observation is that it seems as though  $m = 12$  or  $m = 20$  gives the best filtered signal, with good noise damping in addition to reasonable fidelity close to the edges of the interval. The second is the Gibbs-esque phenomenon occurring at the edges; the DFT assumes a signal is periodic, and the filtered signal reflects this, as it attempts to reconstruct a periodic signal, resulting in the artifacts at the ends of the interval.





