

Exercise session

• Spring 2019 (adapted)

$$\left\{ \begin{array}{l} \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, \quad 0 < x < \pi, \quad t > 0 \end{array} \right.$$

$$u(x, 0) = \sin^2 x \quad (\text{initial condition})$$

$$u(\pi, t) = 0 \quad (\text{Dirichlet BC})$$

$$u_x(0, t) = 0 \quad (\text{Neumann BC})$$

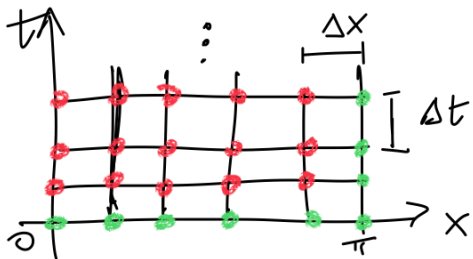
$$\downarrow \frac{\partial u}{\partial x}$$

* Set up the Crank-Nicolson scheme for the numerical solution of this problem.

→ not an exponent!

$$U_i^n \approx u(x_i, t_n)$$

Solution:



* Define Δt and set $t_n = n\Delta t$, $n = 0, 1, 2, \dots$

* Define N and set $\Delta x = \frac{\pi}{N}$ and $x_i = i\Delta x$, $i = 0, 1, 2, \dots, N$.
number of spatial segments

Crank-Nicolson: $\frac{\partial u}{\partial t} = \frac{1}{2} \frac{\partial^2 u}{\partial x^2} + \frac{1}{2} \frac{\partial^2 u}{\partial x^2}$ → at time t_{n+1}

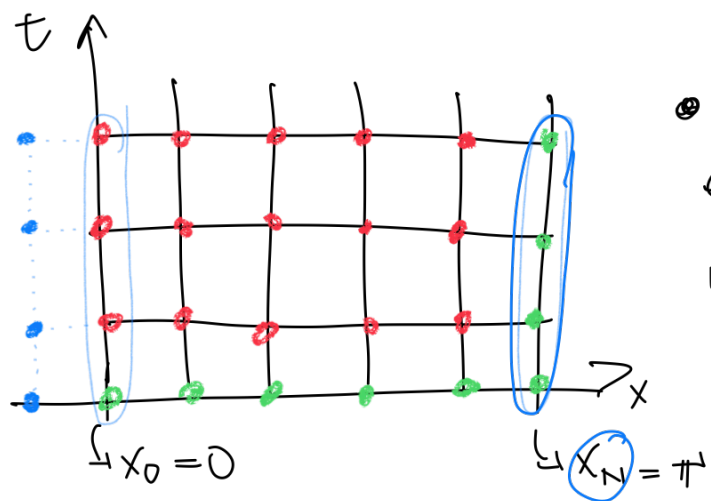
at time t_n

$$\frac{U_i^{n+1} - U_i^n}{\Delta t} = \frac{1}{2} \left(\frac{U_{i-1}^n - 2U_i^n + U_{i+1}^n}{\Delta x^2} \right) + \frac{1}{2} \left(\frac{U_{i-1}^{n+1} - 2U_i^{n+1} + U_{i+1}^{n+1}}{\Delta x^2} \right)$$

- Unknowns: quantities at t_{n+1} :

$$\underbrace{U_i^{n+1}}_{\text{unknown}} - \underbrace{U_i^n}_{\text{known (past)}} = \alpha \underbrace{\left(U_{i-1}^n - 2U_i^n + U_{i+1}^n \right)}_{\text{known (past)}} + \alpha \underbrace{\left(U_{i-1}^{n+1} - 2U_i^{n+1} + U_{i+1}^{n+1} \right)}_{\text{unknown}}, \quad \alpha = \frac{\Delta t}{2\Delta x^2}$$

$$-\alpha U_{i-1}^{n+1} + (1+2\alpha)U_i^{n+1} - \alpha U_{i+1}^{n+1} = \alpha U_{i-1}^n + (1-2\alpha)U_i^n + \alpha U_{i+1}^n \quad | i=0, 1, \dots, N-1$$



• Setting $i=0$ produces a term U_{-1}^{n+1} , which does not exist a priori (not part of the problem). We thus need an equation for U_{-1}^{n+1} . Neumann BC:

$\frac{\partial u}{\partial x} = 0$ for $x=x_0$

→ central discretisation: $\frac{U_1^{n+1} - U_{-1}^{n+1}}{2\Delta x} = 0$

$$\Rightarrow U_{-1}^{n+1} = U_1^{n+1}$$

Dirichlet

= The right (hand side) BC is simpler: $U_N^{n+1} = 0$ for all $n = 0, 1, \dots$

- It remains to handle the initialisation ($n=0$):

$$U_i^0 = \sin^2(x_i) \text{ for } i = 0, 1, \dots, N$$

Final answer:

• $U_i^0 = \sin^2 x_i$, for $i = 0, 1, \dots, N$ (all points in space)

• $U_1^{n+1} - U_{-1}^{n+1} = 0$, for $n = 0, 1, \dots$ (all points in time)

• $-\alpha U_{i-1}^{n+1} + (1+2\alpha)U_i^{n+1} - \alpha U_{i+1}^{n+1} = \alpha U_{i-1}^n + (1-2\alpha)U_i^n + \alpha U_{i+1}^n$, for $\begin{cases} i = 0, 1, \dots, N-1 \\ n = 0, 1, 2, \dots \end{cases}$

• $U_N^{n+1} = 0$ for $n = 0, 1, 2, \dots$ (all points in time)

↳ $U_N^{n+1} \approx u(x_N, t_{n+1}) = u(\pi, t_{n+1}) = 0$

• Autumn 2015 (4D)

$$\frac{\partial^2 u}{\partial x^2} = \frac{\partial u}{\partial t} - 4u, \quad 0 < x < \pi, \quad t > 0$$

non-zero

a) Find all non-trivial solutions of the form $u(x,t) = F(x)G(t)$ that fulfill

the boundary conditions: $\begin{cases} u=0 \text{ for } x=0 \\ \frac{du}{dx}=0 \text{ for } x=\pi \end{cases}$ (for all t)

Solution:

$$\frac{\partial}{\partial x^2} (F(x)G(t)) = \frac{\partial}{\partial t} (F(x)G(t)) - 4F(x)G(t) \Rightarrow \frac{GF''}{FG} = \frac{FG' - 4FG}{FG} \Rightarrow \frac{F''}{F} = \frac{G' - 4G}{G} = K$$

$$\text{ODEs: } \begin{cases} F''(x) = KF(x) \\ G'(t) - 4G(t) = KG(t) \end{cases} \Rightarrow \begin{cases} F'' - KF = 0 \\ G' - (K+4)G = 0 \end{cases}$$

- If $K=0$, then we get $F''=0 \Rightarrow F(x)=Ax+B$, which can only fulfill both BCs if $A=B=0$ (trivial!) \times

- If $K>0$, we get $F(x) = Ae^{\sqrt{K}x} + Be^{-\sqrt{K}x}$ (try yourself to solve that)

→ let's try to plug in the BCs:

$$i) u(0,t) = 0 \Leftrightarrow F(0) = 0 \Leftrightarrow A \cdot 1 + B \cdot 1 = 0 \Leftrightarrow A + B = 0$$

$$ii) \frac{\partial u}{\partial x}(\pi, t) = 0 \Leftrightarrow F'(\pi) = 0 \Leftrightarrow \sqrt{k} A e^{\sqrt{k}\pi} - \sqrt{k} B e^{-\sqrt{k}\pi} = 0 \Leftrightarrow \sqrt{k} e^{\sqrt{k}\pi} (e A - B) = 0$$

Solution:
 $A = B = 0$
 (trivial again!)



- The only possibility left is $k < 0$, that is, $k = -\lambda^2$ ($\lambda \in \mathbb{R}$)

$$\Rightarrow \begin{cases} F''(x) = -\lambda^2 F(x) \\ G'(t) = (4 - \lambda^2) G(t) \end{cases} \Rightarrow \begin{cases} 1) F(x) = A \cos \lambda x + B \sin \lambda x \\ 2) G(t) = C e^{(4 - \lambda^2)t} \end{cases}$$

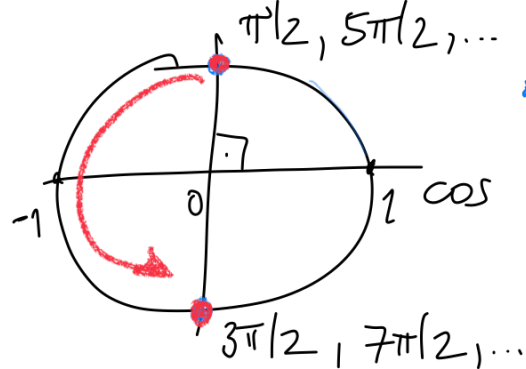
► Boundary conditions:

$$* u(0,t) = 0 \Leftrightarrow F(0) = 0 \Leftrightarrow A \cos 0 + B \sin 0 = 0 \Leftrightarrow \boxed{A = 0}$$

$$* \frac{\partial u}{\partial x} \Big|_{x=\pi} = 0 \Leftrightarrow F'(\pi) = 0 \Leftrightarrow -\lambda A \sin \lambda \pi + \lambda B \cos \lambda \pi = 0 \Leftrightarrow B \cos \lambda \pi = 0$$

trivial!
 \uparrow
 \times
 $\boxed{B = 0}$
 or
 $\boxed{\cos \lambda \pi = 0}$

$$\Rightarrow \cos \lambda \pi = 0$$



* $\cos \lambda \pi = 0$ if $\lambda \pi = m \frac{\pi}{2}$, with m odd

this is the same as writing

$$\boxed{\lambda \pi = (2n-1) \frac{\pi}{2}, \quad n \in \mathbb{N}}$$

$$m = 2n - 1$$

- We therefore have not only one λ , but infinitely many λ_n :

$$\cancel{\lambda} \lambda_n = (2n-1) \frac{\cancel{\pi}}{2} \Rightarrow \boxed{\lambda_n = n - \frac{1}{2}} \quad \parallel \rightarrow \text{for each } \lambda_n, \text{ we have a corresponding pair } F_n(x), G_n(t).$$

$$\rightarrow \text{Solutions: } u_n(x,t) = F_n(x) G_n(t) = \underbrace{[B \sin(\lambda_n x)]}_{F_n(x)} \underbrace{[C e^{(4-\lambda_n^2)t}]}_{G_n(t)} = BC \sin\left(\left(n - \frac{1}{2}\right)x\right) e^{[4 - (n - \frac{1}{2})^2]t}$$

- Since C and B are both arbitrary constants, their product $C \cdot B$ is also arbitrary. We can therefore call $B \cdot C$ as simply an arbitrary constant A_n :

$$u_n(x,t) = A_n \sin\left(\left(n - \frac{1}{2}\right)x\right) e^{[4 - (n - \frac{1}{2})^2]t}, \quad n \in \mathbb{N}$$

b) Find the solution $u(x,t)$ that, in addition to satisfying the PDE and the BCs, also satisfies the IC:

$$u(x,0) = \sin\left(\frac{3x}{2}\right) + 2\sin\left(\frac{5x}{2}\right) + 3\sin\left(\frac{7x}{2}\right)$$

- Superposition: since we have infinitely many solutions $u_n(x,t)$ and this problem admits superposition (why?), then the sum of solutions $u_n(x,t)$ is also a solution: $u(x,t) = \sum_{n=1}^{\infty} A_n \sin\left(\left(n-\frac{1}{2}\right)x\right) e^{[4-(n-1/2)^2]t}$

→ that doesn't mean that all the coefficients A_n will be non-zero!
Maybe so, maybe not (depends on the IC).

$$\text{IC: } u(x,0) = \sin\left(\frac{3x}{2}\right) + 2\sin\left(\frac{5x}{2}\right) + 3\sin\left(\frac{7x}{2}\right) \quad (\text{given})$$

$$u(x,0) = \sum_{n=1}^{\infty} A_n \sin\left(\left(n-\frac{1}{2}\right)x\right) e^0 = \sum_{n=1}^{\infty} A_n \sin\left(\left(n-\frac{1}{2}\right)x\right) \quad (\text{what we have})$$

$$\begin{aligned}
 & \underbrace{A_1}_{=0} \sin\left(\underbrace{\left(1-\frac{1}{2}\right)x}_{1/2}\right) + \underbrace{A_2}_{2} \sin\left(\underbrace{\left(2-\frac{1}{2}\right)x}_{3/2}\right) + \underbrace{A_3}_{3} \sin\left(\underbrace{\left(3-\frac{1}{2}\right)x}_{5/2}\right) + \underbrace{A_4}_{3} \sin\left(\underbrace{\left(4-\frac{1}{2}\right)x}_{7/2}\right) + \dots \\
 & = 1 \cdot \sin\left(\frac{3x}{2}\right) + 2 \sin\left(\frac{5x}{2}\right) + 3 \sin\left(\frac{7x}{2}\right)
 \end{aligned}$$

— Now, let's compare these two: $\left. \begin{array}{l} A_1 = 0 \\ A_2 = 1 \\ A_3 = 2 \\ A_4 = 3 \\ A_5 = 0 \\ A_6 = 0 \\ \vdots \end{array} \right\} \Rightarrow \text{only } n = 2, 3, 4 \text{ contribute!}$

— The solution is therefore:

$$\begin{aligned}
 u(x,t) &= \underbrace{1}_{A_2} \cdot \sin\left(\frac{3x}{2}\right) e^{[4-(3/2)^2]t} + \underbrace{2}_{A_3} \sin\left(\frac{5x}{2}\right) e^{[4-(5/2)^2]t} + \underbrace{3}_{A_4} \sin\left(\frac{7x}{2}\right) e^{[4-(7/2)^2]t} \\
 \Rightarrow u(x,t) &= e^{7t/4} \sin\left(\frac{3x}{2}\right) + 2e^{-9t/4} \sin\left(\frac{5x}{2}\right) + 3e^{-33t/4} \sin\left(\frac{7x}{2}\right) \quad //
 \end{aligned}$$