

Fourier transforms

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(w) e^{iwx} dw$$

$$\hat{f}(w) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(v) e^{-ivw} dv$$

$$\hat{\hat{f}}(w) = f(w)$$

Diff. equation $\xrightarrow{\hat{\cdot}}$ alg. equ. (simpler)
 $\xleftarrow{\hat{\cdot}^{-1}}$

Properties \approx
 $f' \xrightarrow{\quad} (i\omega) \tilde{f}(f)$

But the Fourier
 \Leftarrow transform only
 exists if

$$\int_{-\infty}^{\infty} |f(x)| dx < \infty$$

(very strong)

e.g. e^x is not abs. integrable
 1 is not abs. integrable

In particular, need

$$\lim_{x \rightarrow \pm\infty} f(x) = 0.$$

Heat equation (12.7)

Solve the heat equation

$$\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2}$$

for a bar of infinite
length (very long)

There is no boundary
condition (no boundaries)

Initial condition:

$$u(x, 0) = f(x) \quad \text{for all } x \in \mathbb{R}$$

We can solve this equation using separation of variables

$$u(x, t) = F(x)G(t)$$

$$\text{Get } \begin{cases} F'' + p^2 F = 0 \\ \dot{G} + c^2 p^2 G = 0 \end{cases}$$

$$(k = -p^2)$$

Solutions are of the form

$$F_p(x) = A(p)\cos px + B(p)\sin px$$

$$G_p(t) = e^{-c^2 p^2 t}$$

(No boundary conditions
 $\Rightarrow p$ is any number > 0)

Step 3: In general, none of these particular solutions will solve for the initial conditions.

So we integrate over all p .

$$u(x,t) = \int_0^{\infty} [A(p)\cos px + B(p)\sin px] e^{-cp^2 t} dp$$

①

Using the initial condition:

$$u(x,0) = \int_0^{\infty} [A(p)\cos px + B(p)\sin px] dp$$

= f(x)

(Fourier real integral for F)

$$\Rightarrow A(p) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \cos pv \, dv$$

$$B(p) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \sin pv \, dv$$

$$\textcircled{2} u(x,t) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \left[\int_{-\infty}^{\infty} e^{-c^2 p^2 t} \cos(px - pv) dp \right] dv$$

$$\textcircled{3} u(x,t) = \frac{1}{2c\sqrt{\pi t}} \int_{-\infty}^{\infty} f(v) e^{-\frac{(x-v)^2}{4c^2 t}} dv$$

$$\textcircled{4} = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} f(x + 2cz\sqrt{t}) e^{-z^2} dz$$

Use of Fourier transforms

$$\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2}$$

$u(x, t)$

Apply the Fourier transform on both sides of the equation, thinking of u as a function of x (t is just a parameter)

$\mathcal{F}(u(x, t)) = \text{function of } \omega \text{ \& } t.$
 variable number

$$\tilde{f}(u_t) = c^2 \tilde{f}(u_{xx})$$

taking derivatives
twice with
respect to x

property
2^d

$$= c^2 (i\omega) \tilde{f}(u_x)$$

$$= c^2 (i\omega)(i\omega) \tilde{f}(u)$$

$$= -c^2 \omega^2 \tilde{f}(u)$$

Since \mathcal{F} thinks of u as a function of x

$$\mathcal{F}(u_t) \neq i\omega \mathcal{F}(u)$$

$$\begin{aligned} \mathcal{F}(u_t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u_t e^{-i\omega x} dx \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial u}{\partial t} e^{-i\omega x} dx \\ &= \frac{1}{\sqrt{2\pi}} \frac{\partial}{\partial t} \int_{-\infty}^{\infty} u e^{-i\omega x} dx \end{aligned}$$

$$= \frac{\partial}{\partial t} \tilde{f}(u)$$

$$\Rightarrow \left(\frac{\partial}{\partial t} \tilde{f}(u) = -c^2 \omega^2 \tilde{f}(u) \right)$$

This is an ODE of $\tilde{f}(u)$ with respect to t .
 (ω, t)

A general solution is of the form

$$\tilde{f}(u)(\omega, t) = C(\omega) e^{-c^2 \omega^2 t}$$

$\underbrace{\hspace{10em}}_{t=0 \text{ to find } C(\omega)}$

$$C(\omega) = \tilde{f}(u)(\omega, 0)$$

Recall: $u(x, 0) = f(x)$

$$\Rightarrow \tilde{f}(u)(\omega, 0) = \tilde{f}(f(x))^{(\omega)}$$

$$= \hat{f}(\omega)$$

$$\Rightarrow C(\omega) = \hat{f}(\omega)$$

$$\Rightarrow \tilde{f}(u)(\omega, t) = \hat{f}(\omega) e^{-c^2 \omega^2 t}$$

↑ solution
↑ initial condition

We get u by applying
Inverse FT on both
sides

$$u(x,t) = \mathcal{F}^{-1}(\hat{f}(\omega) e^{-c^2 \omega^2 t})$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{-c^2 \omega^2 t} e^{i\omega x} d\omega$$

$$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(v) e^{i\omega v} dv$$

Note: We have implicitly
assumed that $\mathcal{F}(u)$ exists

$$\Rightarrow \int_{-\infty}^{\infty} |u| dx < \infty.$$

Example: Find a solution $u(x,t)$ to the heat equation with initial condition

$$u(x,0) = \begin{cases} |x| & \text{if } -1 \leq x \leq 1 \\ f(x) & \text{else} \end{cases}$$

We note that $u(x,0)$ is even. Applying the formula (Fourier integral)

$$u(x,t) = \frac{2}{\pi} \int_0^{\infty} \int_0^1 v \cos pv \, dv \cdot (\cos px) e^{-c^2 p^2 t} \, dp$$

\leftarrow $B(p) = 0$
 since f is even

Computing the First integral:

$$u(x,t) = \frac{2}{\pi} \int_0^{\infty} \frac{\cos p x + p \sin p - 1}{p^2} \cdot \rightarrow$$

$$\cos p x e^{-c^2 p^2 t} dp$$

Example:

Use the Fourier transform to solve the PDE

$$t u_x = u_t.$$

satisfying $u(x,0) = f(x)$.

We apply the FT on
both sides (thinking of
 u as a
function of x)

$$\mathcal{F} \left(t \underbrace{u_x^{(x,t)}}_{\substack{\text{only a "number"} \\ \text{in } \mathcal{F}}} \right) = \mathcal{F} \left(u_t^{(x,t)} \right)$$

$$\Rightarrow t \mathcal{F}(u_x) = \mathcal{F}(u_t)$$

$$\Rightarrow t i \omega \mathcal{F}(u) = \frac{\partial}{\partial t} \mathcal{F}(u)$$

This is an ODE of $\mathcal{F}(u)$
with respect to t .

A general solution is

$$\tilde{u}(u)(\omega, t) = C(\omega) e^{t^2/2 i \omega}$$

$$\begin{aligned} C(\omega) &= \tilde{u}(u)(\omega, 0) = \tilde{u}(u(x, 0)) \\ &= \tilde{u}(f(x))_{(\omega)} \\ &= \hat{f}(\omega) \end{aligned}$$

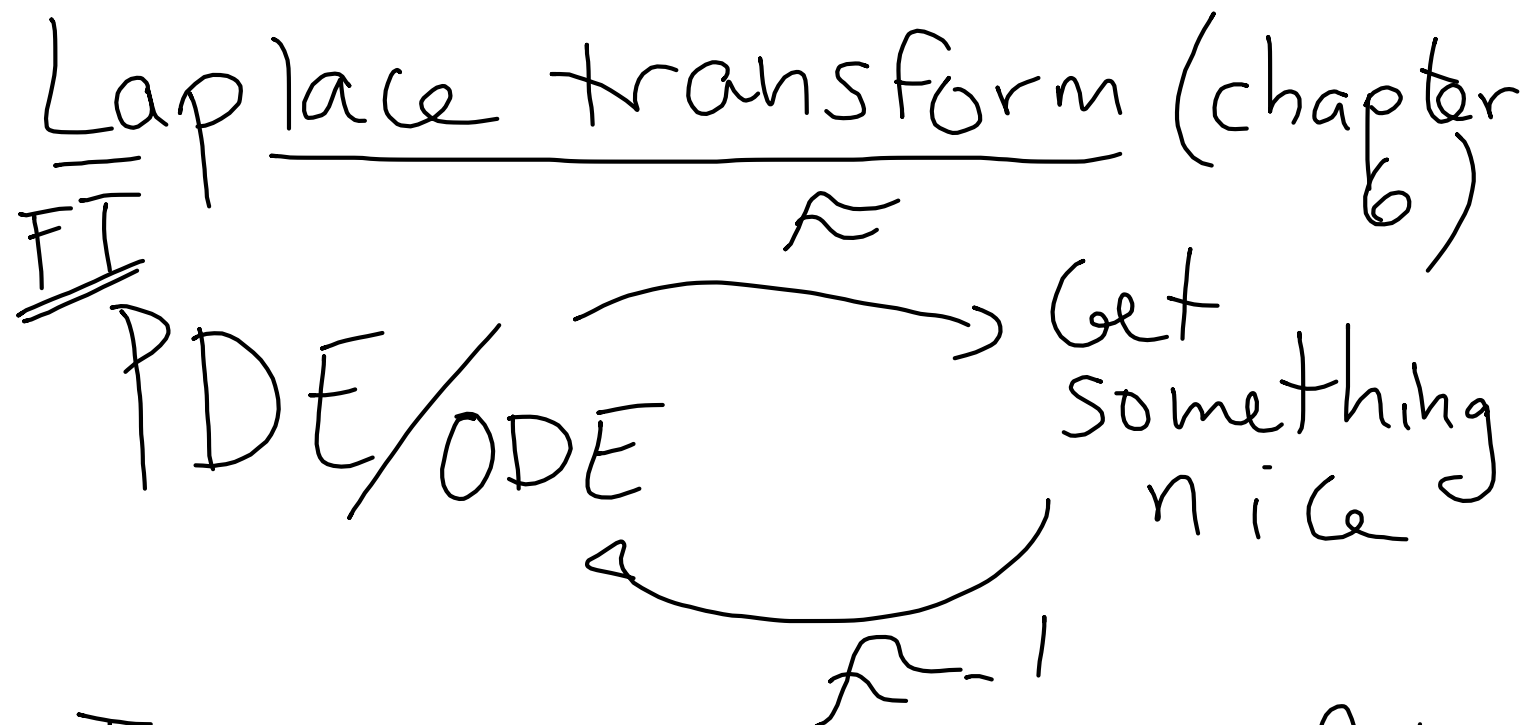
$$\Rightarrow \tilde{u}(u) = \hat{f}(\omega) e^{t^2/2 i \omega}$$

Apply the inverse FT.

$$u(x,t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{t^2/2 i\omega} \cdot e^{i\omega x} d\omega$$
$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{i\omega \left(\frac{t^2}{2} + x\right)} d\omega$$

This is the inverse FT
at $\frac{t^2}{2} + x$

$$= f\left(\frac{t^2}{2} + x\right)$$



In order to apply \mathcal{L} ,
 we assume that \mathcal{L}^{-1}
 $\mathcal{L}(u)$ exists, which
 is not the case often.

Example: $y' = y$

We already know $y = Ce^t$

Say that we want to solve it using FT.

$$\mathcal{F}(y') = \mathcal{F}(y)$$

$$i\omega \mathcal{F}(y) = \mathcal{F}(y)$$

$$\Rightarrow \underbrace{(1 - i\omega)}_{\neq 0} \mathcal{F}(y) = 0$$

$$\Rightarrow \mathcal{F}(y) = 0$$

$$\Rightarrow y = 0$$

The reason why we missed the solution is because

$\mathcal{K}(e^t)$ does not exist.

$$\int_{-\infty}^{\infty} e^t dt = \infty$$

\Rightarrow FT method does not
see this solution.

The goal is to modify
FT to "see"

$$|F(x)| \leq M e^{At}$$

Definition: If $f(t) \ t \geq 0$

Its Laplace transform is
a complex-valued function

$$F(s) = \mathcal{L}(f)(s) = \int_0^{\infty} e^{-st} f(t) dt$$

$$s = \sigma + i\omega$$

$$= \int_0^{\infty} \underbrace{e^{-\sigma t}}_{\text{new}} \cdot \underbrace{e^{-i\omega t} f(t)}_{\text{had before with FT}} dt$$

↳ this will cancel exponential growth in $F(t)$.

In this course,
 S will always be real.