

• Classification

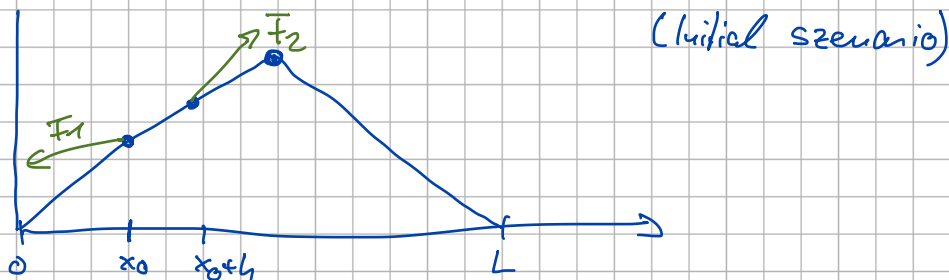
• $\frac{\partial u}{\partial y} - 3 \frac{\partial^2 u}{\partial x^2} + \left(\frac{\partial u}{\partial t}\right)^2 = 0$

- nonlinear
- homogeneous
- order 2

• $\frac{\partial^2 u}{\partial x^2} + \sin(u) - 1 = 0$

- nonlinear
- nonhomogeneous
- order 2

7-7 Derivation of the wave equation



Someone is holding the string (still) in this position and at time $t_0 = 0$ the string is released \leadsto How does the string vibrate?

We have forces acting at every point

Due to assumption 3 (only vertical movement) the horizontal components cancel out

$\underline{F_1} = \begin{pmatrix} -T \\ -T_1 \end{pmatrix}, \quad \underline{F_2} = \begin{pmatrix} T \\ T_2 \end{pmatrix}, \quad \text{i.e. } (\underline{F_1})_x = -T = -(\underline{F_2})_x$

By Newton's second law: any force \underline{F} is mass times acceleration

$\underline{F} = m \cdot \underline{a}, \quad \underline{F}, \underline{a} \in \mathbb{R}^2, \quad u \in \mathbb{R}_+$

So our vertical movement

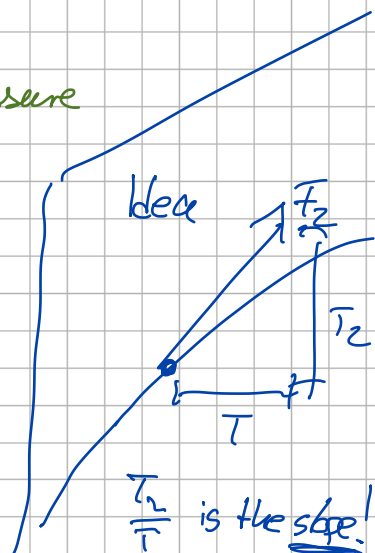
$h \cdot \rho \cdot \frac{\partial^2}{\partial t^2} u(x_0, t) = T_2 - T_1$

length mass per unit length $\left\{ \begin{array}{l} \text{acceleration} \\ \text{force we measure} \end{array} \right.$
 mass

Divide by h : $\rho \cdot \frac{\partial^2}{\partial t^2} u(x_0, t) = \frac{T_2 - T_1}{h}$

Divide by the horizontal (canceling) force T

$\frac{\rho}{T} \frac{\partial^2}{\partial t^2} u(x_0, t) = \frac{\frac{T_2}{T} - \frac{T_1}{T}}{h}$



So $\frac{T_2}{T}$ is the slope of $u(x_0+h, t)$

$\frac{T_1}{T}$ slope at $u(x_0, t)$

But this is $\frac{\partial}{\partial x} u(x_0, t)$

$$\Rightarrow \frac{\rho}{T} \frac{\partial^2}{\partial t^2} u(x_0, t) = \frac{\frac{\partial}{\partial x} u(x_0+h, t) - \frac{\partial}{\partial x} u(x_0, t)}{h}$$

This is a forward difference

Take limit of this equation
 $h \rightarrow 0$

$$\Rightarrow \frac{\rho}{T} \frac{\partial^2}{\partial t^2} u(x_0, t) = \frac{\partial^2}{\partial x^2} u(x_0, t)$$

we introduce $c^2 = \frac{T}{\rho}$ (or $c = \sqrt{\frac{T}{\rho}}$)

$$\Rightarrow \frac{\partial^2}{\partial t^2} u(x, t) = c^2 \frac{\partial^2}{\partial x^2} u(x, t)$$

This is the wave equation.

Classification

- linear
- homogeneous
- order 2

7-9, For the wave equation we obtain

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial^2}{\partial t^2} F(x)G(t) = F(x) \underline{G''(t)}$$

and

$$c^2 \frac{\partial^2 u}{\partial x^2} = c^2 \frac{\partial^2}{\partial x^2} F(x)G(t) = c^2 F''(x)G(t)$$

In our PDE we get

$$\frac{\partial^2 u}{\partial t^2} = \boxed{F(x)G''(t)} = \boxed{c^2 F''(x)G(t)} = c^2 \frac{\partial^2 u}{\partial x^2}$$

Let's rearrange this

$$\frac{G''(t)}{c^2 G(t)} = \frac{F''(x)}{F(x)} = k \quad \text{where } k \in \mathbb{R} \text{ is some constant}$$

$\frac{G''(t)}{c^2 G(t)}$ depends on t
constant in x

$\frac{F''(x)}{F(x)}$ depends on x
constant in t

This means

$$\textcircled{1} \quad \frac{F''(x)}{F(x)} = k \quad \Leftrightarrow F''(x) = kF(x) \quad \Leftrightarrow \boxed{F''(x) - kF(x) = 0}$$

$$\textcircled{2} \quad \frac{G''(t)}{c^2 G(t)} = k \quad \Leftrightarrow \boxed{G''(t) - c^2 k G(t) = 0}$$

These are 2 ODEs (Ordinary differential equations)

We solve these in the following

7-10, The case $k=0$,

$$\text{We get } \textcircled{1} \quad F''(x) - kF(x) = F''(x) = 0$$

$$\Rightarrow F'(x) = A$$

$$\Rightarrow F(x) = Ax + b$$

From the BC we know

$$u(0,t) = F(0)G(t) = 0 \quad \text{for all } t$$

$$u(L,t) = F(L)G(t) = 0 \quad \text{for all } t$$

First solution $G(t) = 0$ for all t

$$u(x,t) = F(x)G(t) = 0 \quad \text{for all } x \text{ and } t$$

a) this is boring

b) the IC $u(x,0) = f(x)$

if $f(x) \neq 0$ this is not a solution.

let's look for other solutions

Second solution $\underline{G(t)} \neq 0$ (at least for one t)

$$\Rightarrow \underline{F(0)} G(t) = 0$$

$$\Rightarrow F(0) = A \cdot 0 + B = 0$$

$$\Leftrightarrow B = 0$$

$$\text{and } F(L) G(t) = 0 \Rightarrow F(L) = A \cdot L + B = 0$$

$$\Leftrightarrow A \cdot L = 0 \Rightarrow A = 0$$

$$\Rightarrow F(x) = 0 \quad \text{for } x \in [0, L]$$

$$\Rightarrow u(x,t) = F(x)G(t) = 0$$

Summary: for $k=0$ only $u(x,t) = 0$ is a solution
this is not of interest.

F=11, The case $k > 0$, let's say $k = \mu^2$, $\mu \in \mathbb{R} \setminus \{0\}$

$$\textcircled{1} F''(x) - k F(x) = 0 \Leftrightarrow \underline{F''(x) = \mu^2 F(x)}$$

$$\textcircled{2} G''(t) - c^2 k G(t) = 0 \Leftrightarrow G''(t) = (c\mu)^2 G(t)$$

For F we take

$$F(x) = e^{\mu x}$$

$$F'(x) = \mu e^{\mu x}$$

$$F''(x) = \mu^2 e^{\mu x}$$

$$F(x) = e^{-\mu x}$$

$$F'(x) = -\mu e^{-\mu x}$$

$$F''(x) = \mu^2 e^{-\mu x}$$

or a combination of both

$$F(x) = A e^{\mu x} + B e^{-\mu x}$$

Similarly

$$G(t) = C e^{\mu c t} + D e^{-\mu c t}$$

(Check!)

(Again!) Check BC.

First solution: if $C=D=0 \Rightarrow G(t)=0$
 (no solution of interest)

$G(t) \neq 0$ for some t .

$$u(0,t) = F(0)G(t) = 0 \Rightarrow F(0) = 0$$

$$u(L,t) = 0 \Rightarrow F(L) = 0$$

$$F(0) = Ae^0 + Be^0 = A+B = 0 \quad (A = -B)$$

$$F(L) = Ae^{\mu L} + Be^{-\mu L} = 0 \quad \left| \begin{array}{l} \text{divide by } e^{-\mu L} \\ \text{or mult. by } e^{\mu L} \end{array} \right.$$

$$\Rightarrow Ae^{2\mu L} + B = 0$$

$$\Leftrightarrow \underbrace{(-1 + e^{2\mu L})}_{\neq 0} A = 0 \Rightarrow A = 0 \Rightarrow B = 0$$

$$\Rightarrow F(x) = 0 \quad x \in [0, L]$$

(no solution of interest)

7-12, (the last) case $k < 0$ or $k = -\lambda^2$ for $\lambda \in \mathbb{R} \setminus \{0\}$

$$\textcircled{1} F''(x) - kF(x) = 0 \Leftrightarrow \underline{F''(x) = kF(x) = -\lambda^2 F(x)}$$

(we look for functions whose 2nd derivative fulfills the eq.)

$$\begin{aligned} F_3(x) &= e^{-i\lambda x} \\ F_3'(x) &= -i\lambda e^{-i\lambda x} \\ F_3''(x) &= -\lambda^2 e^{-i\lambda x} \end{aligned}$$

$$F_1(x) = \sin(\lambda x)$$

$$F_1'(x) = \lambda \cos(\lambda x)$$

$$F_1''(x) = -\lambda^2 \sin(\lambda x) = -\lambda^2 F_1(x)$$

as well as

$$F_2(x) = \cos(\lambda x)$$

$$\Rightarrow F(x) = \underline{A \cos(\lambda x)} + B \sin(\lambda x)$$

(again) BC for $u(x,t) = F(x)G(t)$ ($G(t)$ is not the zero function)

$$\textcircled{1} u(0,t) = 0 \Rightarrow 0 = F(0) = A \cdot 1 + B \cdot 0 = \underline{A = 0}$$

$$\textcircled{2} u(L,t) = 0 \Rightarrow 0 = F(L) = B \sin(\lambda L) = 0$$

we look for zeros
can be nonzero

and $\sin(\lambda L)$ is zero if λL is a multiple of π

So we get $\lambda L = n\pi$, $n=1,2,3,\dots$

$$\Leftrightarrow \lambda = \frac{n\pi}{L}$$

(or our constant k
has to be $k = -\left(\frac{n\pi}{L}\right)^2$)

then $B \cdot \sin\left(\frac{n\pi}{L}\right) = 0$ even for nonzero B !

and for every such λ we get one solution for F :

each $F_n(x) = B \sin\left(\frac{n\pi}{L}x\right)$ solves the first ODE
 \uparrow
set $B=1$

So consider any $k = -\left(\frac{n\pi}{L}\right)^2$ in our second ODE

$$\textcircled{2} G''(t) = -\left(\frac{\pi n}{L}c\right)^2 G(t)$$

$$\text{Define } \lambda_n = \frac{\pi n}{L}c \Rightarrow G''(t) = -\lambda_n^2 G(t)$$

We get a solution for every λ_n as

$$G_n(t) = A_n \cos(\lambda_n t) + B_n \sin(\lambda_n t)$$

With that we finally have solutions

$$u_n(x,t) = F_n(x) \cdot G_n(t) = \left(\underline{A_n \cos(\lambda_n t)} + \underline{B_n \sin(\lambda_n t)} \right) \sin\left(\frac{n\pi}{L}x\right)$$

$$\Rightarrow u(x,t) = \sum_{n=1}^{\infty} u_n(x,t) \quad \text{is a solution}$$

Initial conditions, at $t=0$ $\cos(\lambda_n 0) = 1$
and $\sin(\lambda_n 0) = 0$

$$\underline{f(x)} = u(x,0) = \sum_{n=1}^{\infty} (A_n \cdot 1 + \underline{B_n \cdot 0}) \sin\left(\frac{n\pi}{L}x\right) \quad \text{on } [0,L] \quad x \in [0,L]$$

$$\Rightarrow f(x) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi}{L}x\right) \quad x \in [0,L]$$

use the odd extension of f $f_o(x)$, $x \in [-L,L]$

$$\Rightarrow A_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi}{L}x\right) dx$$

For the B_n we use

$$u(x,t) = (A_n \cos(\omega_n t) + B_n \sin(\omega_n t)) \cdot \sin\left(\frac{n\pi}{L}x\right)$$

$$\frac{\partial}{\partial t} u(x,0) = \sum_{n=1}^{\infty} \left(-A_n \cdot 0 + B_n \frac{c n \pi}{L} \right) \sin\left(\frac{n\pi}{L}x\right)$$
$$= g(x) \quad (\text{initial condition})$$

If we call $d_n = B_n \frac{c n \pi}{L}$

$$g(x) = \sum_{n=1}^{\infty} d_n \sin\left(\frac{n\pi}{L}x\right)$$

use the odd extension $g_0(x) \Rightarrow d_n = \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi}{L}x\right) dx$

$$\Rightarrow B_n = \frac{L}{c n \pi} d_n = \frac{2}{c n \pi} \int_0^L g(x) \sin\left(\frac{n\pi}{L}x\right) dx$$