

Fourier transform

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{i\omega x} d\omega$$

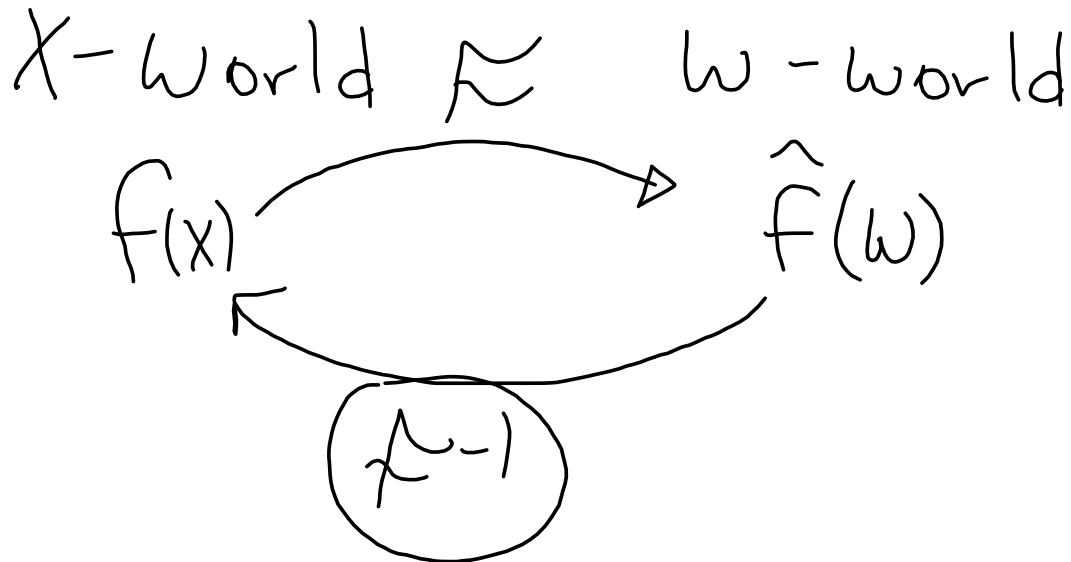
→ Inverse FT of $\hat{f}(\omega)$

$$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(\nu) e^{-i\nu\omega} d\nu$$

Fourier transform

Inverse Fourier transform

$$\mathcal{F}^{-1}(g(\omega)) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(\omega) e^{i\omega x} d\omega$$



Mistake in property 3

$$\mathcal{F}(xf(x))(w) = +i\mathcal{F}'(f(x))(w)$$

Need f to be abs. integrable
 $\int_{-\infty}^{\infty} |f(x)| dx < \infty$ (very strong)

Heat equation: Modeling
very long bars (12.7)

$$\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2}$$

$u(x, t)$: temperature at x
and time $t \geq 0$.

bar of infinite length.

No end points \Rightarrow

No boundary conditions

Initial condition

$$u(x, 0) = f(x) \text{ for all } x \in \mathbb{R}$$

Could solve using the method of separation of variables (3 steps).

Instead, use Fourier transforms

We will only obtain solutions where the FT exists.

We apply FT thinking of u as a function of x (t is considered as a constant)

$$\text{Heat equation: } u_t^{(x,t)} = c^2 u_{xx}^{(x,t)}$$

$$\mathcal{F}(u_t)^{(w,t)} = c^2 \mathcal{F}(u_{xx})^{(w,t)}$$

Cannot use prop. 2 since the derivative is not w.r.t. x .

$$\begin{aligned} &= \text{property 2} \quad c^2 (iw)^2 \mathcal{F}(u) \\ &= -c^2 w^2 \mathcal{F}(u) \end{aligned}$$

$$\begin{aligned}
 \mathcal{F}(u_t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u_t e^{-i\omega x} dx \\
 &= \frac{1}{\sqrt{2\pi}} \frac{\partial}{\partial t} \int_{-\infty}^{\infty} u e^{-i\omega x} dx \\
 &= \frac{\partial}{\partial t} \mathcal{F}(u)
 \end{aligned}$$

$$\Rightarrow \frac{\partial}{\partial t} \mathcal{F}(u)^{(\omega, t)} = -c^2 \omega^2 \mathcal{F}(u)^{(\omega, t)}$$

This is an ODE of $\mathcal{F}(u)$ with respect to t . (ω as a constant)

Solution to ODE

$$\tilde{u}(u)(\omega, t) = \underbrace{C(\omega)}_? e^{-c^2 \omega^2 t}$$

Using the initial condition

$$\begin{aligned} C(\omega) &= \tilde{u}(u)(\omega, 0) \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \underbrace{u(x, 0)}_{\text{initial condition}} e^{-i\omega x} dx \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\omega x} dx \\ &= \hat{f}(\omega) \end{aligned}$$

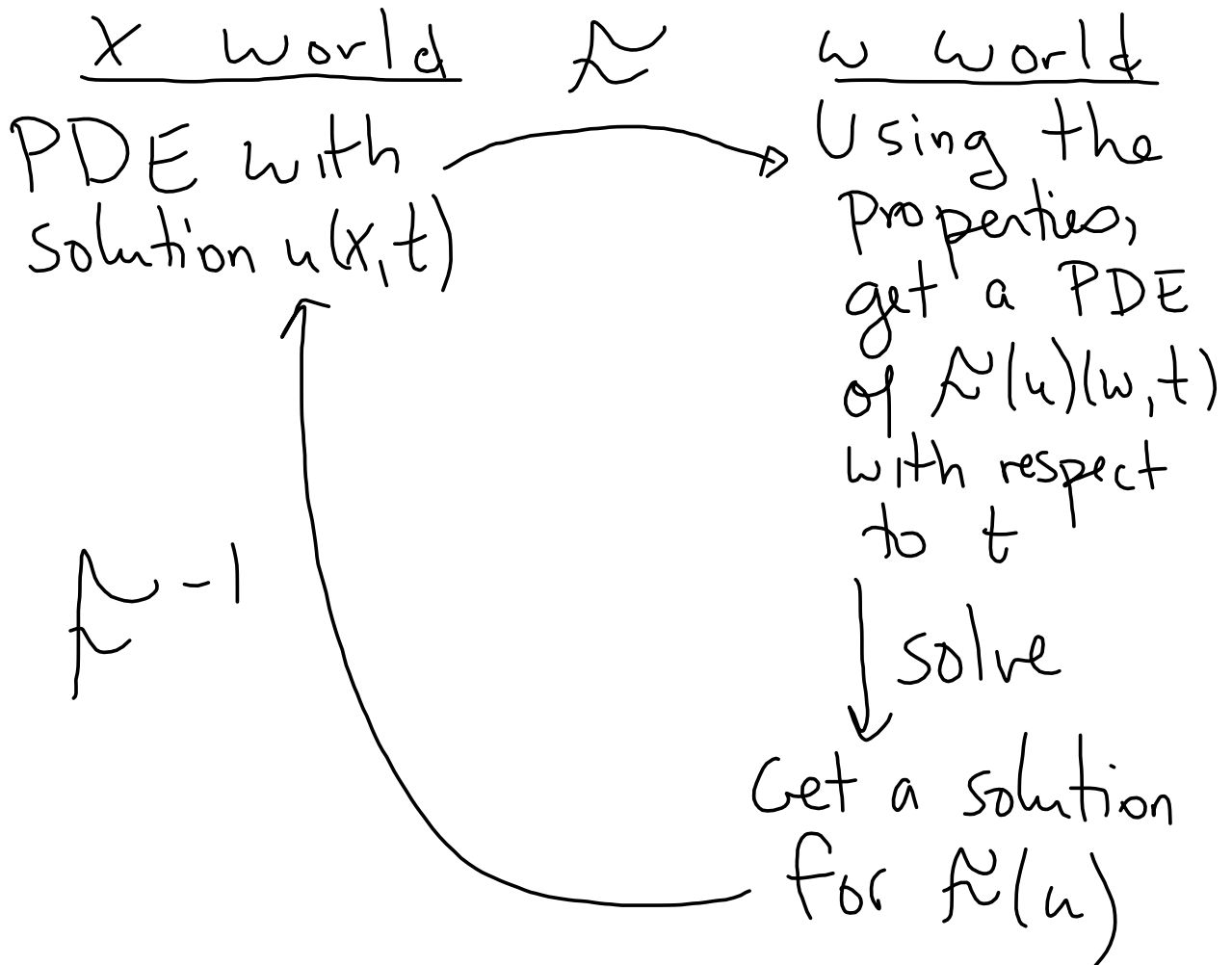
$$\Rightarrow \boxed{\tilde{u}(w, t) = \hat{f}(w) e^{-c^2 w^2 t}}$$

To find $u(x, t)$, we apply the inverse FT.

$$\boxed{u(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(w) e^{-c^2 w^2 t} e^{iwx} dw}$$

where

$$\hat{f}(w) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(v) e^{-i w v} dv$$



Example Find the solution to the heat equation with initial conditions

$$f(x) = u(x, 0) = \begin{cases} |x| & \text{if } -1 \leq x \leq 1 \\ 0 & \text{else} \end{cases}$$

f is even.

$$u(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(w) e^{-c^2 w^2 t} e^{iwx} dw$$

$$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(v) e^{-i\omega v} dv$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-1}^1 |v| e^{-i\omega v} dv$$

even \rightsquigarrow cosine integral

$$= \frac{2}{\sqrt{2\pi}} \int_0^1 v \cos(\omega v) dv$$

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$$\begin{aligned} \Rightarrow u(x,t) &= \frac{2}{2\pi} \int_{-\infty}^{\infty} \int_0^1 v \cos(\omega v) dv e^{-c^2 \omega^2 t} e^{i\omega x} d\omega \\ &= \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{\cos \omega + \omega \sin \omega - 1}{\omega^2} e^{-c^2 \omega^2 t} e^{i\omega x} d\omega \end{aligned}$$

Example: Use the FT to solve the PDE

$$t u_x = u_t$$

Find $u(x,t)$

Apply FT on both sides
with respect to x .

$$\mathcal{F}(t u_x) = \mathcal{F}(u_t)$$

$$\Rightarrow \underbrace{t \mathcal{F}(u_x)}_{\substack{\text{Property} \\ (2)}} = \underbrace{\mathcal{F}(u_t)}_{\substack{\text{can take } \partial/\partial t \\ \text{out}}}$$

$$\Rightarrow t \cdot i \omega \mathcal{F}(u) = \frac{\partial}{\partial t} \mathcal{F}(u)$$

\leadsto ODE of $\mathcal{F}(u)$ with
respect to t . (ω is a
constant)

$$\Rightarrow \mathcal{K}(u)(\omega, t) = C(\omega) e^{\frac{t^2}{2} i\omega}$$

$$C(\omega) = \mathcal{K}(u)(\omega, 0) = \hat{f}(\omega).$$

Applying the inverse
FT

$$\begin{aligned} \Rightarrow u(x, t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{i\omega \frac{t^2}{2}} e^{i\omega x} d\omega \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{i\omega \left(\frac{t^2}{2} + x\right)} d\omega \end{aligned}$$

This is the inverse FT
of \hat{f} evaluated at $\frac{t^2}{2} + x$

$$= \boxed{f\left(\frac{t^2}{2} + x\right)}$$

Different form of the
solution to the heat equation
using convolution

$$u(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(w) e^{-c^2 w^2 t} e^{iwx} dw$$

$$= \int_{-\infty}^{\infty} \hat{f}(w) \hat{g}(w) e^{iwx} dw, \text{ where}$$

$$\hat{g}(w) = \frac{1}{\sqrt{2\pi}} e^{-c^2 w^2 t}$$

This is the inverse FT of a product of FTs.

$$= (f \star g)(x, t)$$

(t is a constant)
initial condition

We Find $g(x, t)$.

Already know from previous example

$$\mathcal{F}(e^{-ax^2}) = \frac{1}{\sqrt{2a}} e^{-\omega^2/4a}$$

$a \neq 0$.

We will modify $\hat{g}(\omega)$ so that it looks like

$$\text{Choose } a = \frac{1}{4c^2 t}$$

(to get the same exponent)

$$\Rightarrow \frac{1}{\sqrt{2a}} = \sqrt{2c^2 t}$$

$$\hat{g}(\omega) = \frac{1}{\sqrt{2\pi}} e^{-c^2 \omega^2 t}$$

$$= \frac{1}{\sqrt{2\pi}} \cdot \frac{1}{\sqrt{2c^2 t}} \cdot \sqrt{2c^2 t} e^{-c^2 \omega^2 t}$$

$$= \frac{1}{\sqrt{2\pi}} \cdot \frac{1}{\sqrt{2c^2 t}} \cdot \boxed{\frac{1}{\sqrt{2a}} e^{-\omega^2 / 4a}}$$

where $a = \frac{1}{4c^2 t}$.

Applying inverse FT

$$\Rightarrow g(x, t) = \frac{1}{\sqrt{2\pi}} \cdot \frac{1}{\sqrt{2c^2 t}} e^{-ax^2}$$

$$= \frac{1}{2c\sqrt{\pi t}} e^{-\frac{x^2}{4c^2 t}}$$

Using the definition of convolution:

$$u(x, t) = (f * g) = \int_{-\infty}^{\infty} f(p) g(x-p) dp$$

$$\left[= \frac{1}{2c\sqrt{\pi t}} \int_{-\infty}^{\infty} f(p) e^{-\frac{(x-p)^2}{4c^2 t}} dp \right]$$