

# Løysingsforslag for TMA4120, Øving 5

October 3, 2016

**11.7.2)** La  $f : \mathbb{R} \rightarrow \mathbb{R}$  vere gitt av

$$f(x) = \begin{cases} \frac{\pi}{2} \sin x & \text{for } |x| \leq \pi \\ 0 & \text{ellers.} \end{cases}$$

Sidan  $f$  er odde, er  $A(\omega) = 0$  for alle  $\omega \geq 0$ . Vi reknar ut

$$\begin{aligned} B(\omega) &= \frac{1}{\pi} \int_{-\infty}^{\infty} f(x) \sin \omega x dx = \frac{1}{2} \int_{-\pi}^{\pi} \sin x \sin \omega x dx \\ &= \frac{1}{2} \left[ \frac{\cos x \sin \omega x - \omega \sin x \cos \omega x}{\omega^2 - 1} \right]_{x=-\pi}^{\pi} \\ &= \frac{\sin \pi \omega}{1 - \omega^2}. \end{aligned}$$

Det er klart at  $f$  oppfyller krava for representasjonsteoremet for Fourierintegral side 513. Sjekk at du ser kvifor! Sidan  $f$  ikkje har diskontinuitetspunkt, er middelverdien overalt lik funksjonsverdien sjølv. Derfor er

$$\int_0^{\infty} (A(\omega) \cos \omega x + B(\omega) \sin \omega x) d\omega = \int_0^{\infty} \frac{\sin \pi \omega}{1 - \omega^2} \sin \omega x d\omega = f(x) = \begin{cases} \frac{\pi}{2} \sin x & \text{for } |x| \leq \pi \\ 0 & \text{ellers.} \end{cases}$$

Det var dette vi skulle vise.

**11.7.11)** La  $g : \mathbb{R} \rightarrow \mathbb{R}$  vere gitt ved  $g(x) = f(|x|)$ . Sidan  $g$  tilfredsstillar krava i Teorem 1 og er like, er  $g(x)$ , og dermed  $f(x)$ , gitt ved (10) på side 515. Reknar ut  $A(w)$ :

$$\begin{aligned} A(w) &= \frac{2}{\pi} \int_0^{\infty} g(v) \cos wv dv \\ &= \frac{2}{\pi} \int_0^{\pi} \sin v \cos wv dv \\ &= \frac{2}{\pi} \left[ \frac{w \sin v \sin wv + \cos v \cos wv}{w^2 - 1} \right]_0^{\pi} \\ &= \frac{2 \cos w\pi + 2}{\pi(1 - w^2)}. \end{aligned}$$

Vi får dermed

$$f(x) = \int_0^\infty \frac{2 \cos w\pi + 2}{\pi(1-w^2)} \cos wx dw$$

**11.7.19)** Vi bruker same metode som over og latar som om  $f$  er ein del av ein odde funksjon definert på heile  $\mathbb{R}$ . Dermed er  $f$  gitt ved (11), der

$$\begin{aligned} B(w) &= \frac{2}{\pi} \int_0^\infty f(v) \sin wv dv \\ &= \frac{2}{\pi} \int_0^1 e^v \sin wv dv \\ &= \frac{2}{\pi} \left( [e^v \sin wv]_0^1 - \int_0^1 w e^v \cos wv dv \right) \\ &= \frac{2}{\pi} \left( e \sin w - [e^v w \cos wv]_0^1 - \int_0^1 w^2 e^v \sin wv dv \right) \\ &= \frac{2}{\pi} \left( e \sin w + w - ew \cos w - w^2 \int_0^1 e^v \sin wv dv \right) \\ \Rightarrow B(w) &= \frac{2}{\pi} \cdot \frac{w + e(\sin w - w \cos w)}{1 + w^2}. \end{aligned}$$

Det gir

$$f(x) = \int_0^\infty \frac{2(w + e(\sin w - w \cos w))}{\pi(1+w^2)} \sin wx dw$$

**11.9.4)**

$$\begin{aligned} \sqrt{2\pi} \hat{f}(w) &= \int_{-\infty}^\infty f(x) e^{-iwx} dx \\ &= \int_{-\infty}^0 e^{kx} e^{-iwx} dx \\ &= \int_{-\infty}^0 e^{(k-iw)x} dx \\ &= \left[ \frac{1}{k-iw} e^{(k-iw)x} \right]_{x=-\infty}^0 \\ &= \frac{1}{k-iw} \\ \Rightarrow \hat{f}(w) &= \frac{1}{\sqrt{2\pi}(k-iw)} \end{aligned}$$

**11.9.9)**

$$f(x) = \begin{cases} |x| & \text{for } -1 < x < 1 \\ 0 & \text{ellers} \end{cases}$$

$$\begin{aligned}
\hat{f}(w) &= \frac{1}{\sqrt{2\pi}} \left( \int_{-1}^0 -xe^{-iw x} dx + \int_0^1 xe^{-iw x} dx \right) \\
&= \frac{1}{\sqrt{2\pi}} \left( \frac{1}{iw} xe^{-iw x} \Big|_{-1}^0 - \int_{-1}^0 \frac{1}{iw} e^{-iw x} dx - \frac{1}{iw} xe^{-iw x} \Big|_0^1 + \int_0^1 \frac{1}{iw} e^{-iw x} dx \right) \\
&= \frac{1}{\sqrt{2\pi}} \left( \frac{1}{iw} e^{iw} + \frac{1}{(iw)^2} e^{-iw x} \Big|_{-1}^0 - \frac{1}{iw} e^{-iw} - \frac{1}{(iw)^2} e^{-iw x} \Big|_0^1 \right) \\
&= \frac{1}{\sqrt{2\pi}} \left( \frac{1}{iw} e^{iw} - \frac{1}{w^2} + \frac{1}{w^2} e^{iw} - \frac{1}{iw} e^{-iw} + \frac{1}{w^2} e^{-iw} - \frac{1}{w^2} \right) \\
&= \frac{1}{\sqrt{2\pi}} \left( -\frac{2}{w^2} + \frac{2}{w^2} w \sin w + \frac{2}{w^2} \cos w \right) \\
&= \frac{\sqrt{2}}{\sqrt{\pi} w^2} (\cos w + w \sin w - 1)
\end{aligned}$$

Her brukte vi at  $e^{iw} = \cos w + i \sin w$ . Det går også an å utnytte at  $f$  er like og skrive  $\hat{f}(w) = \hat{f}_c(w) = \sqrt{\frac{2}{\pi}} \int_0^1 x \cos wx dx$ , der  $\hat{f}_c(w)$  er Fourier-cosinus-transformen av  $f$  fra 11.8.

**H)** Vi har

$$\begin{aligned}
\hat{f}(w) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-iw x} dx \\
&= \frac{1}{\sqrt{2\pi}} \left( - \int_{-\infty}^0 e^x e^{-iw x} dx + \int_0^{\infty} e^{-x} e^{-iw x} dx \right) \\
&= \frac{1}{\sqrt{2\pi}} \left( - \int_{-\infty}^0 e^{(1-iw)x} dx + \int_0^{\infty} e^{-(1+iw)x} dx \right) \\
&= \frac{1}{\sqrt{2\pi}} \left( - \left[ \frac{1}{1-iw} e^{(1-iw)x} \right]_{-\infty}^0 - \left[ \frac{1}{1+iw} e^{-(1+iw)x} \right]_0^{\infty} \right) \\
&= \frac{1}{\sqrt{2\pi}} \left( \frac{1}{1+iw} - \frac{1}{1-iw} \right) \\
&= -\frac{\sqrt{2}iw}{\sqrt{\pi}(1+w^2)}.
\end{aligned}$$

Dette gir

$$\begin{aligned}
f(x) &= \mathcal{F}^{-1} \left( -\frac{\sqrt{2}iw}{\sqrt{\pi}(1+w^2)} \right) \\
&= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} -\frac{\sqrt{2}iwe^{iw x}}{\sqrt{\pi}(1+w^2)} dw \\
\Rightarrow f(1) &= e^{-1} = -\frac{1}{\pi} \int_{-\infty}^{\infty} \frac{iwe^{iw}}{(1+w^2)} dw.
\end{aligned}$$

(Her kan vi også skrive  $f(x) = \frac{2}{\pi} \int_0^\infty \frac{w \sin wx}{1+w^2} dw$ , siden  $f$  er odde.) Sidan

$$\begin{aligned} I &:= \int_0^\infty \frac{w \sin w}{1+w^2} dw = \frac{1}{2} \int_{-\infty}^\infty \frac{w \sin w}{1+w^2} dw \\ &= -\frac{1}{2} \operatorname{Re} \int_{-\infty}^\infty \frac{iwe^{iw}}{1+w^2} dw, \end{aligned}$$

er

$$I = -\frac{1}{2} \operatorname{Re} \left( -\frac{\pi}{e} \right) = \frac{\pi}{2e}.$$

**I)**

$$\begin{aligned} \hat{f}(w) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^\infty \cos t e^{-t^2} e^{-iwt} dt \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^\infty \frac{1}{2} (e^{it} + e^{-it}) e^{-t^2} e^{-iwt} dt \\ &= \frac{1}{2\sqrt{2\pi}} \int_{-\infty}^\infty e^{-t^2} (e^{-i(w-1)t} + e^{-(w+1)t}) dt \\ &= \frac{1}{2} (\mathcal{F}(e^{-t^2})(w-1) + \mathcal{F}(e^{-t^2})(w+1)) \\ &= \frac{1}{2\sqrt{2}} \left( e^{-(w-1)^2/4} + e^{-(w+1)^2/4} \right) \end{aligned}$$

som også kan skrivast  $\frac{1}{\sqrt{2}} e^{-(w^2+1)/4} \cosh(w/2)$ . Her har vi brukt 9 i Tabell III side 536.

**J)** Vi bruker formelen

$$\mathcal{F}(f * g) = \sqrt{2\pi} \mathcal{F}(f) \mathcal{F}(g)$$

og får

$$\begin{aligned} \mathcal{F}(h) &= \mathcal{F}(e^{-x^2} * e^{-x^2}) = \sqrt{2\pi} \mathcal{F}(e^{-x^2}) \mathcal{F}(e^{-x^2}) \\ &= \sqrt{2\pi} \left( \frac{1}{\sqrt{2}} e^{-w^2/4} \right)^2 \\ &= \sqrt{\frac{\pi}{2}} e^{-w^2/2}. \end{aligned}$$

Vi ser at med  $a = \frac{1}{2}$  i formel 9, Table III, har vi

$$\begin{aligned} \mathcal{F}(e^{-x^2/2}) &= e^{-w^2/2} = \sqrt{\frac{2}{\pi}} \mathcal{F}(h) \\ \Rightarrow h(x) &= \sqrt{\frac{\pi}{2}} e^{-x^2/2}. \end{aligned}$$

**K)** Vi vil ta Fouriertransform av begge sider, og vi byrjar med å rekne ut følgande:

$$\begin{aligned}
 \mathcal{F}(e^{-a|x|}) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-a|x|} e^{-iwx} dx \\
 &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^0 e^{(a-iw)x} dx + \frac{1}{\sqrt{2\pi}} \int_0^{\infty} e^{-(a+iw)x} dx \\
 &= \frac{1}{\sqrt{2\pi}} \left( \frac{1}{a-iw} + \frac{1}{a+iw} \right) \\
 &= \frac{2a}{\sqrt{2\pi}(a^2+w^2)}.
 \end{aligned}$$

Integralet i likninga er konvolusjonen  $f * e^{-3|x|}$ , så, sidan  $\mathcal{F}(f * g) = \sqrt{2\pi} \mathcal{F}(f) \mathcal{F}(g)$ , kan vi Fouriertransformere begge sider av likninga og få

$$\begin{aligned}
 \hat{f}(w) - \sqrt{2\pi} \frac{6}{\sqrt{2\pi}(9+w^2)} \hat{f}(w) &= \frac{6}{\sqrt{2\pi}(9+w^2)} \\
 \Rightarrow \frac{3+w^2}{9+w^2} \hat{f}(w) &= \frac{6}{\sqrt{2\pi}(9+w^2)} \\
 \Rightarrow \hat{f}(w) &= \frac{6}{\sqrt{2\pi}(3+w^2)}.
 \end{aligned}$$

Med  $a = \sqrt{3}$  over gir dette

$$f(x) = \sqrt{3} e^{-\sqrt{3}|x|}.$$

**La)**

$$\begin{aligned}
 \hat{f}(w) &= \frac{1}{\sqrt{2\pi}} \int_{-1}^1 e^{-iwx} dx \\
 &= \frac{1}{\sqrt{2\pi}} \left[ \frac{-e^{-iwx}}{iw} \right]_{-1}^1 \\
 &= \frac{-e^{-iw} + e^{iw}}{iw\sqrt{2\pi}} \\
 &= \frac{2i \sin w}{iw\sqrt{2\pi}} \\
 &= \sqrt{\frac{2}{\pi}} \frac{\sin w}{w}
 \end{aligned}$$

$$\begin{aligned}
\hat{g}(w) &= \frac{1}{\sqrt{2\pi}} \int_0^\infty e^{-(1+iw)x} dx \\
&= \frac{1}{\sqrt{2\pi}} \left[ \frac{-e^{-(1+iw)x}}{1+iw} \right]_0^\infty \\
&= \frac{1}{\sqrt{2\pi}} \frac{1}{1+iw} \\
&= \frac{1}{\sqrt{2\pi}} \frac{1-iw}{1+w^2}
\end{aligned}$$

**Lb)** Vi har  $\mathcal{F}(h) = \mathcal{F}(f * g) = \sqrt{2\pi} \mathcal{F}(f) \mathcal{F}(g)$ . Vi tek invers Fouriertransform og får

$$\begin{aligned}
h(x) &= \mathcal{F}^{-1}(\sqrt{2\pi} \hat{f} \hat{g}) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^\infty \sqrt{2\pi} \hat{f}(w) \hat{g}(w) e^{iwx} dw \\
&= \frac{1}{\pi} \int_{-\infty}^\infty \frac{(1-iw) \sin w}{w(1+w^2)} e^{iwx} dw.
\end{aligned}$$

Set vi  $x = 0$ , får vi

$$\begin{aligned}
\frac{1}{\pi} \int_{-\infty}^\infty \frac{(1-iw) \sin w}{w(1+w^2)} dw &= h(0) \\
&= \int_{-\infty}^\infty f(-p)g(p) dp \\
&= \int_{-1}^1 g(p) dp \\
&= \int_0^1 e^{-p} dp \\
&= 1 - e^{-1}
\end{aligned}$$

Integralet vi skal finne er realdelen av integralet på venstresida. Sidan høgresida er reell, får vi

$$\frac{1}{\pi} \int_{-\infty}^\infty \frac{\sin w}{w(1+w^2)} dw = 1 - e^{-1},$$

altså

$$\int_{-\infty}^\infty \frac{\sin w}{w(1+w^2)} dw = \pi(1 - e^{-1}).$$

**M)** Vi har  $\sin w = i(e^{iw} - e^{-iw})/2$ . Sidan  $f$  er den same funksjonen som i L), har vi

$$\begin{aligned}
\hat{f}(w) &= \sqrt{\frac{2}{\pi}} \frac{\sin w}{w} \\
&= \frac{i}{\sqrt{2\pi}} \frac{e^{iw} - e^{-iw}}{w}.
\end{aligned}$$

Så finn vi  $\mathcal{F}(f * f)$ :

$$\begin{aligned}\mathcal{F}(f * f) &= \sqrt{2\pi} \hat{f} \cdot \hat{f} = \frac{-1}{\sqrt{2\pi}} \left( \frac{e^{iw} - e^{-iw}}{w} \right)^2 \\ &= \frac{-1}{\sqrt{2\pi}} \frac{e^{2iw} - 2 + e^{-2iw}}{w^2}.\end{aligned}$$

Sidan  $f * f$  er kontinuerleg, gir invers Fouriertransform

$$(f * f)(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{-1}{\sqrt{2\pi}} \frac{e^{2iw} - 2 + e^{-2iw}}{w^2} e^{iwx} dw.$$

Set vi  $x = 3$ , får vi

$$-\frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{5iw} - 2e^{3iw} + e^{iw}}{w^2} dw = (f * f)(3) = 0,$$

sidan  $(f * f)(3) = \int_{-\infty}^{\infty} f(p)f(3-p)dp = \int_{-1}^1 f(3-p)dp = \int_2^4 f(t)dt = 0$ .  
Integralet vi er ute etter er, sidan  $e^{iaw} = \cos aw + i \sin aw$ , realdelen av integralet på venstresida. Dermed er

$$\int_{-\infty}^{\infty} \frac{\cos 5w - 2 \cos 3w + \cos w}{w^2} dw = 0.$$