# THE POROUS MEDIA EQUATION

$$u_{\perp} = \Delta u_{m}$$
 m>1 (PME)

The PME is a nonlinear diffusion equation:

Note that the diffusion coefficient mu<sup>m-1</sup> vanishes when u=0, making the equation degenerate. Therefore, the PME is sometimes called the slow diffusion equation (SDE).

In contrast, if ormal, the diffusion coefficient goes to so when u > 0.

This is <u>fast</u> diffusion, which we shall not consider have.

If m=1, we just get the standard heat equation. From now on, we assume m>1.

#### MODELLING

Let u be the density of gas in a porous medium: Then

$$u_t + div(uv) = 0$$

where u is the flow role. Davey's law:

And the gas law:  $p = CU^{\gamma}$  ( $\gamma = 1$  for an ideal gas).

Thus we get

ut = ck div (upu) = cky div (upu) = cky Dull
The constent is easily scaled away.

The equation also models the expansion of the fiveball in a nuclear explotion. The 'porons" medium" is the air. It does not have time to move. Heat transfer is by radiation, move efficient at higher temporative. (I have not studied the details.)

SELF-SIMILAR SOLUTIONS: THE BARGUBLATT SOLUTION Look for solutions on the form  $u(t,x) = t^{-\alpha} f(xt^{-\beta})$ . Conservation of mass  $\Leftrightarrow x = \beta d$ . Also  $\alpha = \frac{d}{(m-1)d+2}$ 

Bavenblatt:  $u(t,x) = \left(A - \frac{B(m-1)}{2m} \frac{x^2}{t^p}\right)_{t}^{m-1}$ 

# THE GENERALISED PME

(PME) only makes sense for u>0.

But we can extend it:  $u_t = \Delta(1u1^M \text{sgn u})$ .

This way, the diffusion coefficient minimized is nonnegative for all u.

(A negative diffusion coefficient leads to be havior like the backward heat equation, which is not well behaved.)

More generally, we consider

$$u_{\xi} = \Delta(\varphi(u))$$
 (gPME)

where  $\varphi$  is a strictly increasing function 12-312 with  $\varphi(0) = 0$ . We shall always assume that  $\varphi$  is locally Lipschitz. By Rademacher's theorem, then  $\varphi$  is differentiable almost everywhere. (In one dimension, this is easy, since a Lipschitz function is clearly absolutely continuous.)

A fancier way of saying this is  $\varphi \in W_{loc}(12)$ .

## Some PROTERTIES OF THE GPME

This section is all about CLASSICAL solutions which vanish, along with relevant derivatives, as IXI-20. We ship some technical defails.

Prop Let u and v be solutions of (gPME). Then  $\int_{\mathbb{R}^d} (u-v)_+ dx$  is a (non-strictly) decreasing function of t.

Proof sketch Let w = Q(u) - Q(v), so  $(u - v)_{\xi} = \Delta w$ . Let  $p: \mathbb{R} \to \mathbb{R}$  be smooth with  $p' \ge 0$  and  $0 \le p \le 1$ . They

$$\int_{\mathbb{R}^d} p(w) (u-v)_{\xi} dx = \int_{\mathbb{R}^d} p(w) \Delta w dx$$

$$= -\int_{\mathbb{R}^d} p(w) |\nabla w|^2 dx \leq 0$$

\* assuming reasonably fast decay at infinity, so the boundary term vanishes

Now let  $p(w) \rightarrow [w>o]$  (Ivanson bracket). Because  $\varphi$  is strictly increasing, [w>o] = [u-v>o]and  $[u-v>o](u-v)_t = \partial_t (u-v)_+$ , leading to  $\partial_t \int_{\mathbb{R}^d} (u-v)_+ dx \leq 0$ 

#### COROLLEY (MONOTONICITY)

Let u, v be solutions of (gPME) with initial data  $u_0, v_0$ . If  $u_0 \le v_0$  then  $u \le v$ .

Pq. Sa(u-v)+ dx =0 for t=0, so it remains zero I

### COROLLARY (MIXIMON PRINCIPLE)

If a = u0 = b then a = u = b. Also, ||u||20 = |u0||20
Pf. The constant functions are solutions.

COROLLARY (L' CONTRACTION) || u(t)-v(t)||, & ||40-v<sub>0</sub>||, Pf. Both  $\int (u-v)_+ dx$  end  $\int (v-u)_+ dx$  ene decreasing.

Add the two together.

Mass Conservation: Judx = Judx

Pf  $\frac{\partial}{\partial t} \int_{\mathbb{R}^d} u \, dx = \int_{\mathbb{R}^d} u_t \, dx = \int_{\mathbb{R}^d} \varphi(u) \, dx = 0$ 

## CONSTRUCTING & SOLUTION

Assume given the INITIAL VALUE PROBLEM:

$$u_t = \Delta(\phi \circ u) \times cR^{\lambda}, t>0$$

$$u(0, \times) = u_0(\times)$$

We assume that of is increasing and locally Lipschitz, and up & L'(Rd) of Locally Lipschitz, and up & L'(Rd) of Locally Lipschitz, where we solution up Locally of Locally Lipschitz, and up & L'(Rd) of Charles of L'(Rd) of the lipschitz, and locally Lip

$$\int_{0}^{T} \int_{\mathbb{R}^{d}} (u \psi_{t} + \varphi(u) \underline{A} \psi) dx dt + \int_{\mathbb{R}^{d}} u_{0}(x) 4(0,x) dx = 0$$

for all  $CPED((-10,T)\times\mathbb{R}^d)$ . For SIMPLICITY, Put d=1. Everything still works the same for d>1.

We choose a time step k and look for approximations  $u^n(x) \neq u(nk,x)$  for all k.

We DO NOT DISCRETIZE IN SPACE, but we still employ the discrete Leplacian, and put

$$\frac{k}{n_{M1}(x) - n_{M1}(x)} = \nabla^{\mu}(clon) := \frac{\mu_{M1}(x+\mu) - 5d_{M1}(x) + b_{M1}(x-\mu)}{\mu_{M2}}$$

where we write  $q^n(x) = q(u^n(x))$  for brevity.

Introduce  $\lambda = \frac{k}{h^2}$ , and write the scheme as

$$u^{N+1}(x) = u^{N}(x) - 2\lambda cp(u^{N}(x))$$
  
+  $\lambda (q(u^{N}(x+h)) + q(u^{N}(x-h)))$ 

We want the scheme to be monotone, meaning the RHS should be an increasing function of  $u^{n}(x)$ ,  $u^{n}(x+h)$  and  $u^{n}(x-h)$ . The latter two are easy since of is increasing. The first requires  $2\lambda K \leq 1$ , where  $K = \sup Q'(s)$ .

To be precise, we should take the supremum

Then un will take values in this

We Sould use ess sup here...

interval for all n.

Our upper limit on  $\lambda$  can be written as  $k \leq \frac{h^2}{2K}$ .

NOTATION Write our scheme on the form  $u^{n+1} = G[u^n]$ .

So long as  $a \le u \le v \le b$   $(u, v \in L^{\infty}(R^n))$ ,

elso  $G[a] \le G[u] \le G[b]$ .

From this monotonicity, we get desirable proporties of our scheme, resembling proporties of (gPME).

LEMM Jource) dx is independent of x.

Proof. Trivial, because & Dh(Qou) dx =0

LERMA Let u, or be solutions of our scheme. Then  $\int (u^n-v^n)_+ dx$  is a (non-strictly) decreasing function of n.

Proof We went to prove  $\int_{\mathbb{R}} (G[u] - G[v])_{+} dx \leq \int_{\mathbb{R}} (u - v)_{+} dx$ 

Consider the pointwise maximum uvu.

Clearly  $u \leq uvu$ , so  $G[u] \leq G[uvu]$ .

Therefore  $G[u] - G[v] \leq G[uvu] - G[v]$ But then also  $(G[u] - G[v])_{+} \leq G[uvu] - G[v]$ Since  $PHS \geq 0$  everywhere. Therefore  $(G[u] - G[v])_{+} dx \leq G(G[uvu] - G[v]) dx$ 

 $\int (G[u]-G[v])_{+} dx \leq \int (G[uvv]-G[v]) dx$   $= \int (G[uvv]-gv]) dx$   $= \int (u-v)_{+} dx$ 

COROLLARY Our scheme is L'-contractive: 1.0., llundi-untill, < llun-until.

Translation: Write  $\tau_S u(x) = u(x-\delta)$ .

Our scheme is translation invarient, so  $(\tau_S u^N)$  is a solution of the scheme with

initial value  $\tau_S u^0$ . From L' contraction we get  $\|u^N - \tau_S u^N\|_1 \le \|u^0 - \tau_S u^0\|_1$ Put  $\omega(\varepsilon) = \sup \|u^0 - \tau_S u^0\|_1$ .  $\|S\| \le \varepsilon$ 

It is well known that  $\lim_{\epsilon \to 0} \omega(\epsilon) = 0$ , and we have just seen that  $\|u^n - \tau_S u^n\|_1 \le \omega(\epsilon)$  for dl n and  $(\delta | \le \epsilon)$ .

Tightness: Let  $\chi \in C^{\infty}(\mathbb{R})$  with  $\chi(x) = 0$  for  $|x| < \frac{1}{\epsilon}$ ,  $\chi(x) = 1$  for |x| > 1, end  $0 \le \chi(x) \le 1$  for all x.

Let  $\chi_{\mathbb{R}}(x) = \chi(x/\mathbb{R})$   $\frac{R_2}{\epsilon}$ 

Then

 $\int \chi_{R} u^{nH} dx - \int \chi_{R} u^{n} dx = k \int \chi_{R} \Delta_{h} (Q \circ u^{n}) dx = k \int (\chi_{Q} - 1) \Delta_{h} (Q \circ u^{n}) dx$   $= k \int \Delta_{h} \chi_{R} (Q \circ u^{n}) dx = k O(R^{-2}) \text{ as } R \rightarrow \infty.$ 

By induction, SXRUMAX - SXRUOdx = kn O(RZ) = O(TRZ).

# L' compressens:

By Kolmogorov, the set {un | new, kn < T} is precompact in LICIR).

### EQUICONTINUITY out TIME:

We need to estimate 11 unti-un11,.

To that end, introduce a standard molifier p, and  $g_{\epsilon}(x) = \epsilon^{-1} p(x/\epsilon)$  as usual. We find

To estimate  $\|\Delta_h g_{\varepsilon}\|_1$ , consider any  $\psi \in C_{\varepsilon}^2(\mathbb{R})$ :  $\psi(x+h) - \psi(x) = \int_0^h \psi'(x+t) dt$ , so  $\psi(x+h) - 2\psi(x) + \psi(x-h) = \int_0^h (\psi'(x+t) - \psi'(x-h+t)) dt$   $= \int_0^h \int_0^o \psi''(x+s+t) ds dt \text{ and so}$ 

112/411 = 1-5 215 1 2/4 "(x+x+f) as af 1 dx

= 6 -2 5 -4 1 14" (x+s+t) | dx ds dt = 114" 11,

and so

 $\|\Delta_{h} \rho_{\epsilon}\|_{1} \leq \|\rho_{\epsilon}^{"}\|_{1} = \int_{\mathbb{R}^{2}} e^{-3} \rho^{"}(x/\epsilon) dx = e^{-2} \|\rho^{"}\|_{1}$ 

and so

Write 
$$u=u^n$$
. The  $u(x)-p_e*u(x)=\int_{-\epsilon}^{\epsilon}(u(x)-u(x-y))p_{\epsilon}(y)dy$ 
and therefore

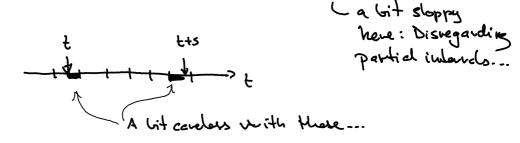
S

$$\|u^{n+j}-u^{n}\|_{1} \leq \overline{\omega}(jk)$$
  
where  $\overline{\omega}(s) = \inf_{\varepsilon > 0} \left(2\omega(\varepsilon) + \frac{sC}{\varepsilon^{2}}\right)$ 

Lemma

Pf. Given  $\eta > 0$ , pick  $\varepsilon$  so that  $\omega(\varepsilon) < \gamma$ . Then if  $S < C\varepsilon^2 \eta$ , we get  $\overline{\omega}(S) \in 2\omega(\varepsilon) + \frac{SC}{\varepsilon^2} < 3\eta$ .  $\square$  Define an approximate adultion  $\tilde{u}(t,x)$  by linear inderpolation: With  $t^{N} = nk$ , put  $\tilde{u} = \frac{t^{N+1}-t}{k} u^{N} + \frac{t-t^{N}}{k} u^{N+1}$  for  $t \in Ct^{N}$ ,  $t^{M+1}$ .

Then we get  $\|\tilde{u}(t+s) - \tilde{u}(t)\|_{1} \leq \tilde{w}(s)$ 



We can now call upon the Arzelà-Ascoli theorem: Writing our approximate solution  $\tilde{u}$  as  $\tilde{u}(h_ik): [0,T] \rightarrow L^1(\mathbb{R}),$  there is a sequence  $(h_i,h_i) \rightarrow (0,0)$  with  $k_i \leq \frac{h_i^2}{2K}$  which converges uniformly to some  $u \in C(C_iT); L^1(IR))$  Since q is Lipschitz (combat K), e(so  $\|\varphi(\tilde{u}(h,k))-\varphi(u)\|_1 \leq K \|\tilde{u}(h,k)-u\|_1 \rightarrow 0$  uniformly.

The following calculation will show that the limit u is a weak solution.

Put th=nk for n=0,1,..., while nk≤T. Then

$$\int_{\mathbb{R}} \left( \tilde{u} \, \psi_{t} + \varphi(\tilde{u}) \Delta \psi \right) dx dt$$

$$= \sum_{N \geqslant 0} \int_{t^{N} \tilde{R}}^{t^{N+1}} \left( \tilde{u} \, \psi_{t} + \varphi(\tilde{u}) \Delta_{h} \psi \right) dx dt$$

$$= \sum_{N \geqslant 0} \int_{\mathbb{R}} \left( u^{N+1} \psi(t^{N+1}) - u^{N} \psi(t^{N}) \right) dx$$

$$+ \sum_{N \geqslant 0} \int_{\mathbb{R}} \left( -\tilde{u}_{t} \psi + \varphi(u^{N}) \Delta_{h} \psi \right) dx dt$$

$$= - \int_{\mathbb{R}} u^{0} \psi(0) dx$$

$$+ \sum_{N \geqslant 0} \int_{t^{N}} \int_{\mathbb{R}} \left( -\frac{u^{N+1} - u^{N}}{k} + \Delta_{h} \varphi(u^{N}) \right) \psi dx dt$$

$$= - \int_{\mathbb{R}} u^{0} \psi(0) dx$$

$$= - \int_{\mathbb{R}} u^{0} \psi(0) dx$$

In this calculation,  $\cong$  indicates that the difference between the two sides goes to zero when  $(h,k) \rightarrow (0,0)$ .

Justification on the next page!

In the first ~ on the previous pege, we have neglected

which goes to zero because |\sh4-&4|las ->0 as h ->0, while ||q(\widelex)||, remains bounded.

In the second =, we neglicited

 We have proved Existence:

Assume 9: R > R is strictly increasing and locally Lipschitz,

Then the initial value problem  $u_t = \Delta(dou) \times cR^d, t>0$   $u(0,x) = u_0(x)$ 

has a solution

ue La([o,T]×Rd) nC([o,T]; L'(Rd)), in the week sonse:

 $\int_{\mathbb{R}^{d}} \left( u \psi_{t} + \varphi(u) \underline{u} \psi \right) dx dt + \int_{\mathbb{R}^{d}} u_{0}(x) 4(0,x) dx = 0$ 

for all QED ((-10,T)×Rd).

Assuming we can prove that only one such solution exists, i.e., uniqueness, then we must get the same limit for every choice of  $(h_i, k_i) \rightarrow (0,0)$  (satisfying  $k_i \in \frac{h_i^2}{2K}$ ), and so we conclude that  $\tilde{u}(h_i k) \rightarrow u$  when  $(h_i, k) \rightarrow (0,0)$  with  $k \leq \frac{h^2}{2K}$ .